

Long Baseline Neutrino Physics

**Manfred Lindner
Max-Planck-Institute, Heidelberg**



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1. Introduction

The Birth of the Neutrino



Before 1930: neutron \rightarrow proton + e^-
2-body decay \rightarrow monoenergetic spectrum expected

experiment: continuous β -decay spectrum

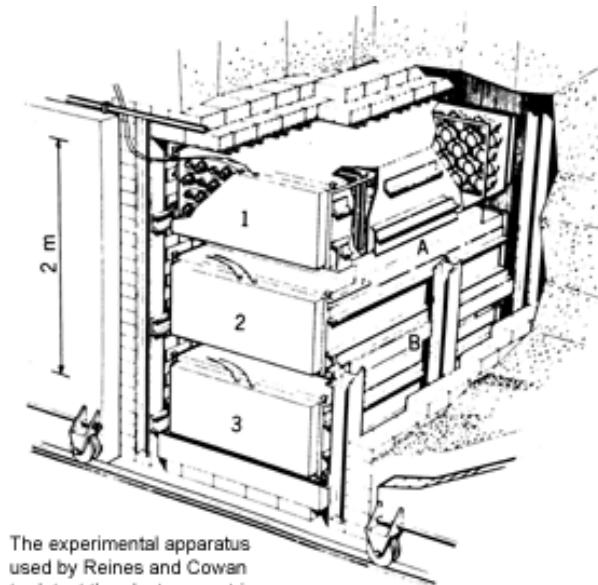
Pauli: energy-momentum conservation

- postulate new particle
- invisible, since $Q=0$
- spin $1/2$, ...

Letter to Tübingen Dec. 1930 ...
... will never be detected

•

- Cowen & Reines 1954-56 project "poltergeist"
 - detection of reactor neutrinos
 - Nobel price for F. Reines 1995

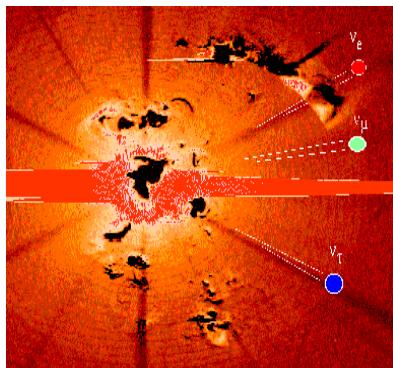


The experimental apparatus used by Reines and Cowan to detect the electron neutrino.

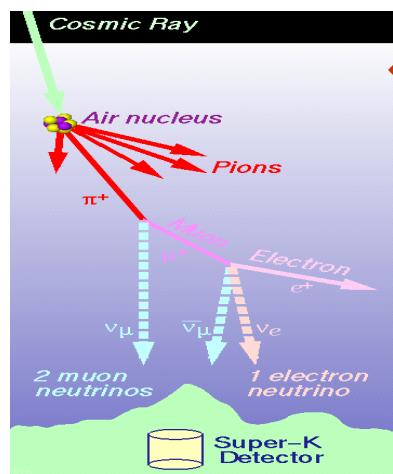
Neutrino Sources & New Physics



← Sun



← Cosmology



← Atmosphere



← Earth

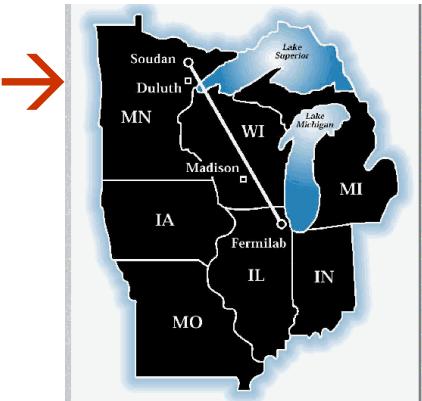
Astronomy: →
Supernovae
UHE ν's

...

Reactors →



Accelerators →
Laboratories



Physics Beyond the Standard Model

Theoretical arguments:

SM does not exist without cutoff

Higgs-doublet = only simplest extension

Gauge hierarchy problem

Why: 3 generations , fermion representations

Many parameters (9+? Masses, 4+? Mixings)

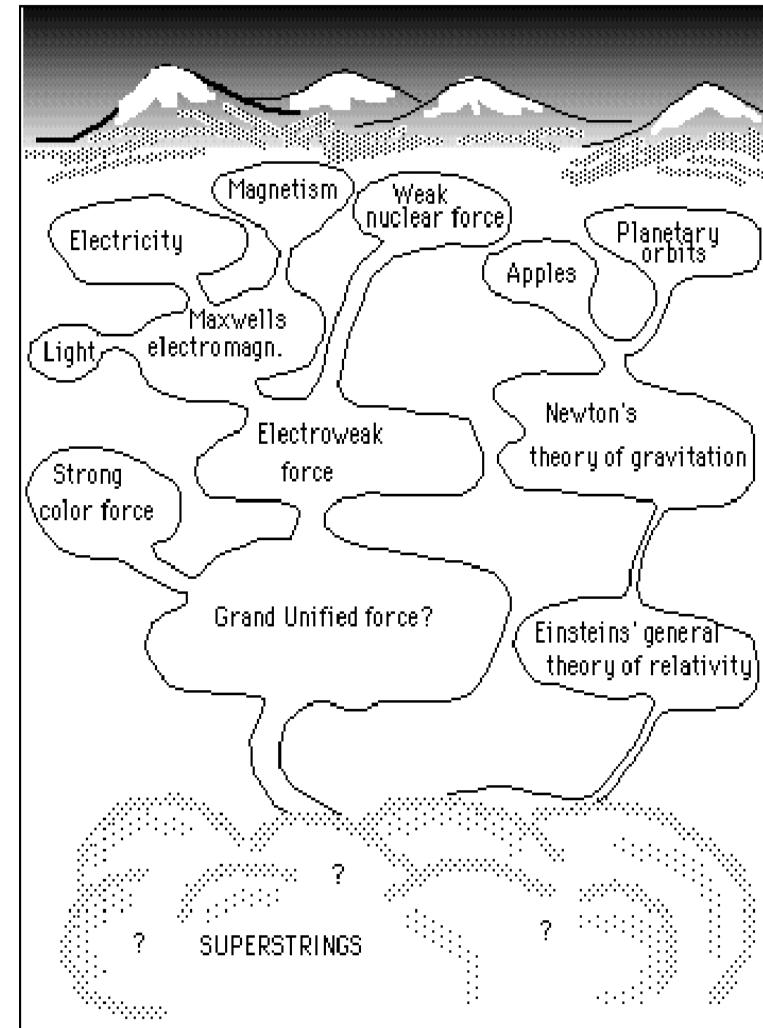
Charge quantisation, unification: GUTs, ...,

Gravitation, ...

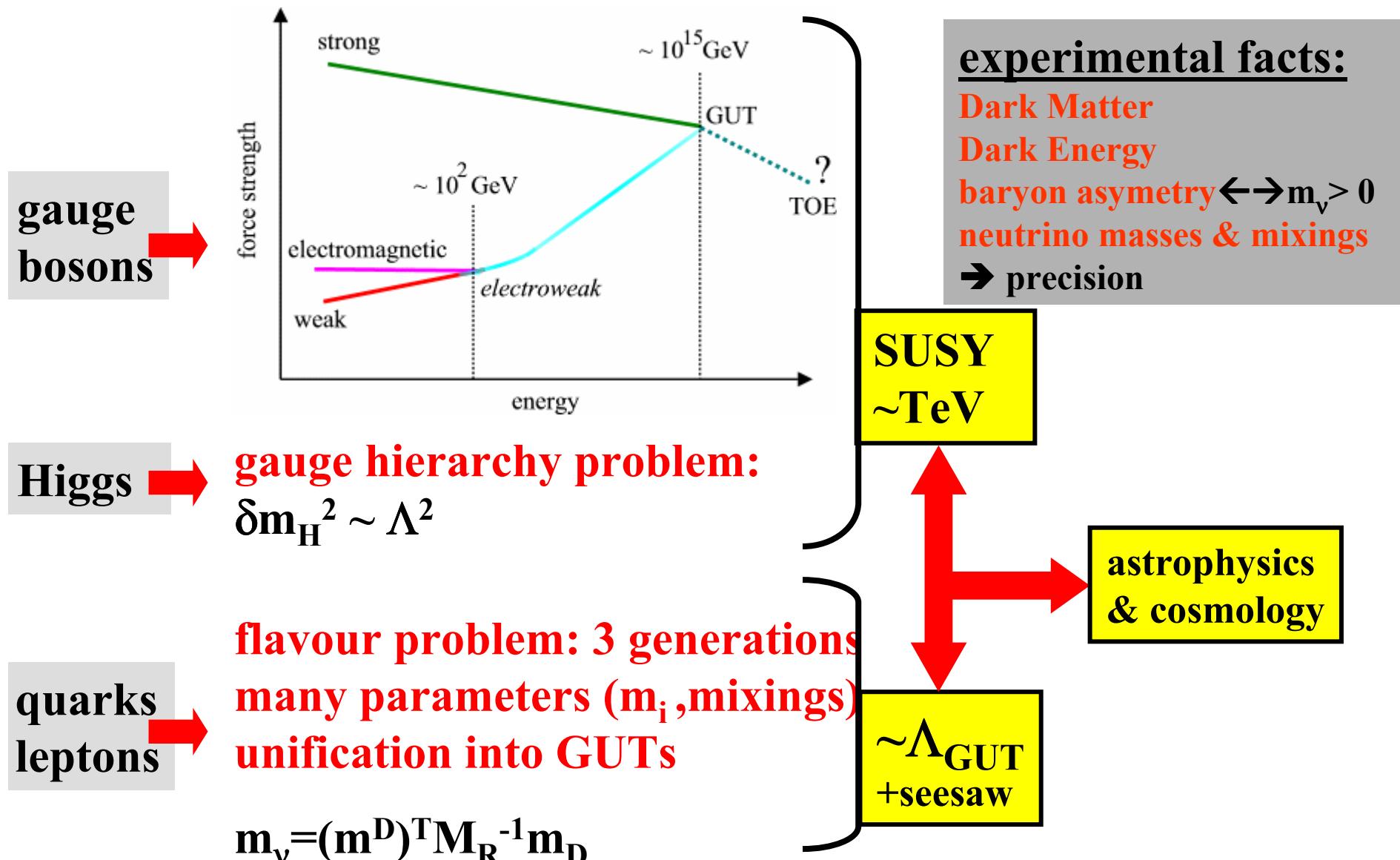
2 directions: Sym. breaking & Flavour

Experimental facts:

- Dark Matter & Dark Energy exist!
- Neutrino masses have been detected!
- Baryon asymmetry of the universe $\leftrightarrow m_\nu > 0$
- physics beyond the standard model
- results \leftrightarrow implications for theory



New Physics Beyond the SM



2. Neutrino Masses

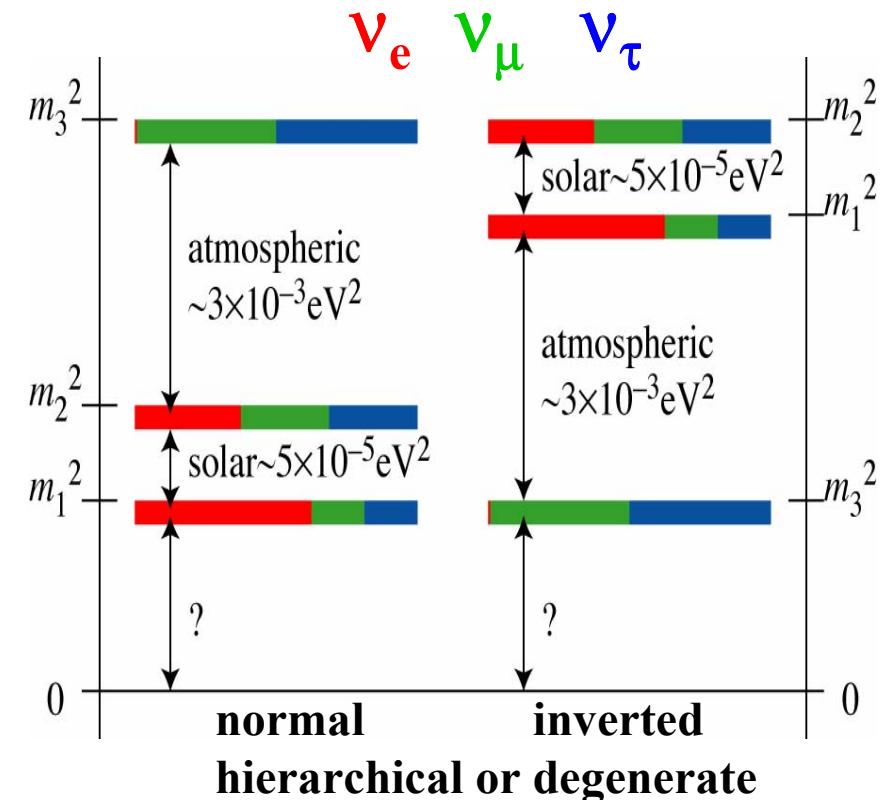
Parameters for 3 Light Neutrinos

mass & mixing parameters: m_1 , Δm^2_{21} , $|\Delta m^2_{31}|$, $\text{sign}(\Delta m^2_{31})$

$$U = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix} \text{diag}(e^{i\alpha}, e^{i\beta}, 1)$$

questions:

- Dirac or Majorana
- absolute mass scale: m_1
- mass ordering: $\text{sgn}(\Delta m^2_{31})$
- how small is θ_{13} , θ_{23} maximal?
- leptonic CP violation
- LSND ↔ sterile neutrino(s)
- L/E pattern of oscillations



4 Ways to measure Masses & Mixings

4 different methods:

- kinematical
- lepton number violation \leftrightarrow Majorana
- astrophysics & cosmology
- oscillations

4 Ways to measure Masses & Mixings

4 different methods:

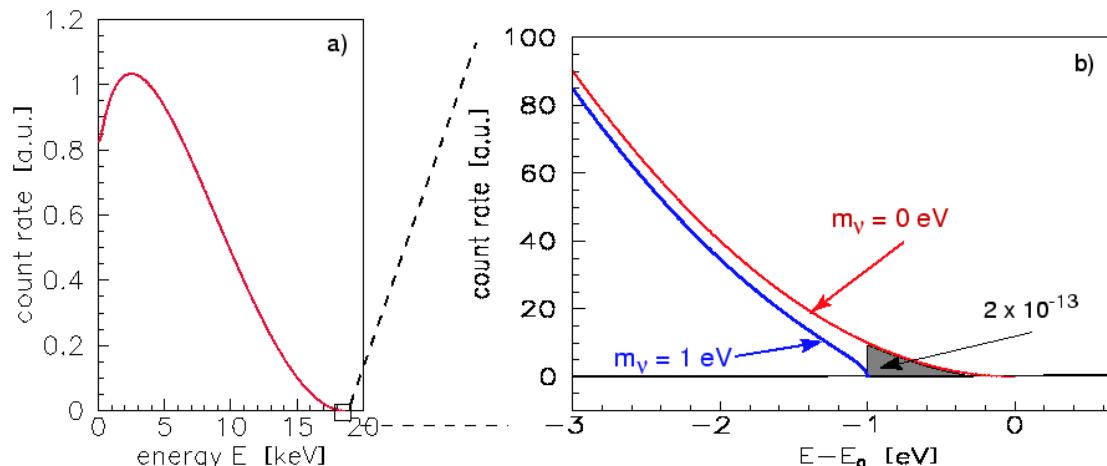
- **kinematical**
- lepton number violation \leftrightarrow Majorana
- astrophysics & cosmology
- oscillations

Summary: Kinematical Mass Determination

Relativistic kinematics:

$$E^2 = p^2 + m^2; \quad \sum p_i^\mu = \sum p_f^\mu$$

Endpoint of decays:



Bounds:

“Elektron-Neutrino”: $m < 2.2$ eV (Mainz, Troitsk)

“Muon-Neutrino”: $m < 170$ keV

“Tau-Neutrino”: $m < 15.5$ MeV

Sensitivity \Leftrightarrow degenerate ν -spectrum

\Rightarrow **Oscillations:** $\Delta m_{ij}^2 \ll m_i^2 \Rightarrow \sum m_i^2 |U_{ei}|^2 < (2.2 \text{ eV})^2$

Future: KATRIN: $\longrightarrow 0.25$ eV

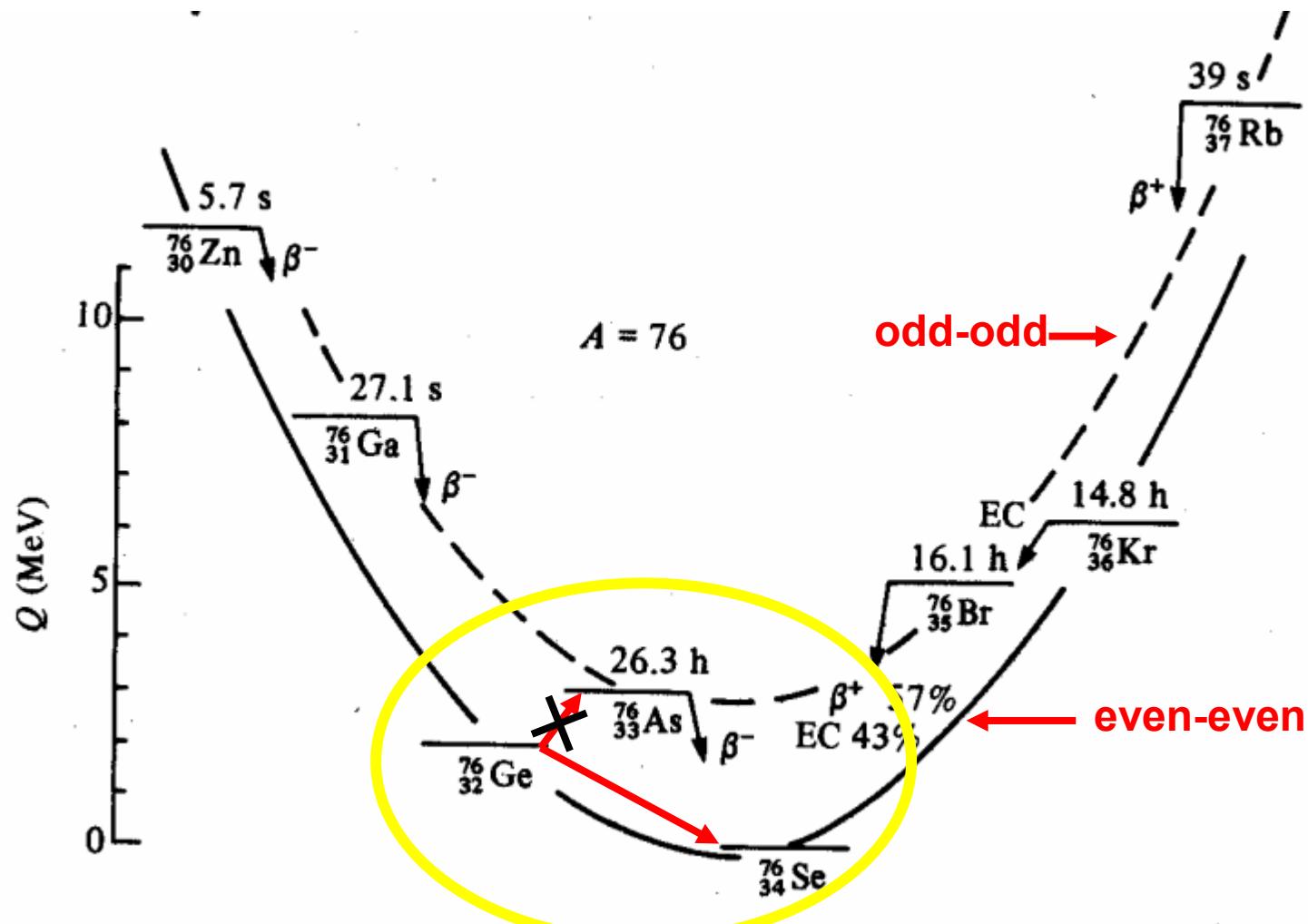
\Leftrightarrow cosmological bound
(WMAP; 2df galaxy survey, Lyman α , ...)

4 Ways to measure Masses & Mixings

4 different methods:

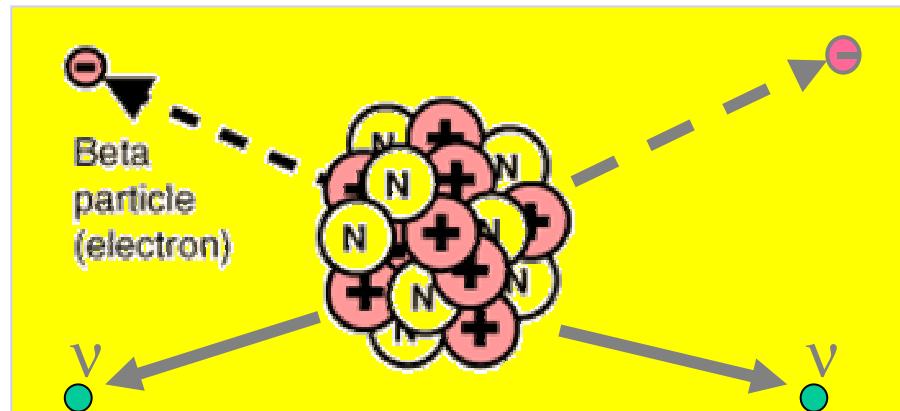
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Mass Parabolas

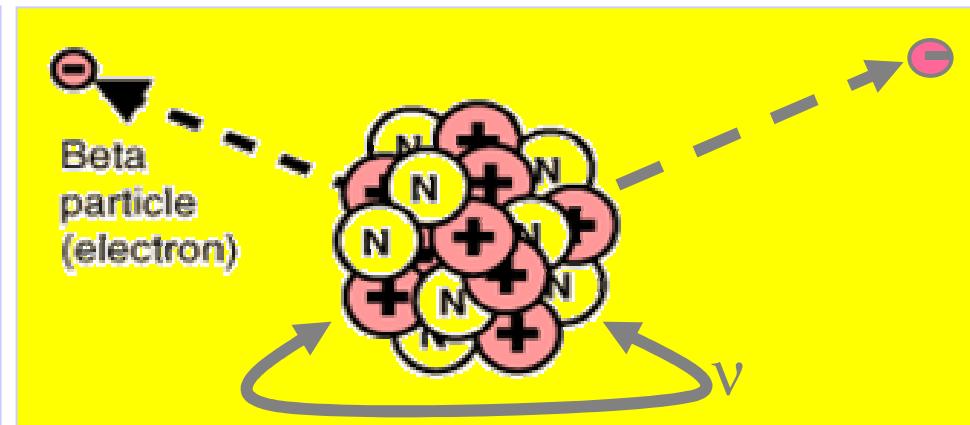


Ground states of even-even nuclei: 0^+

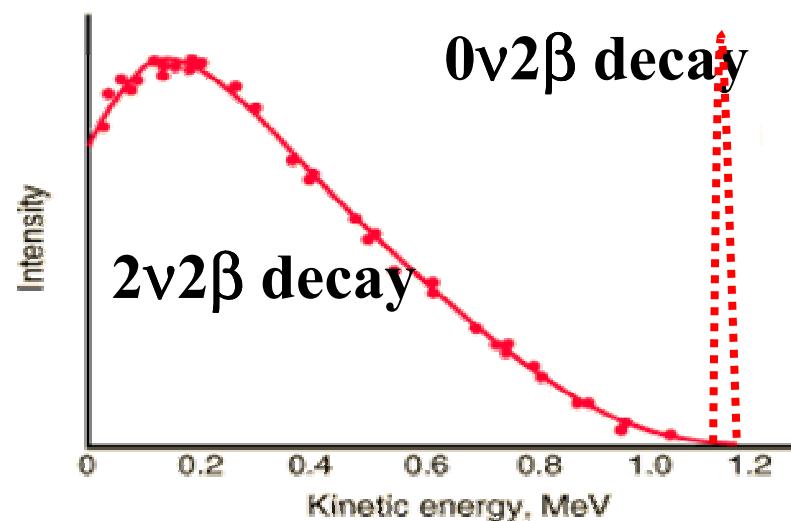
Lepton Number Violation & $0\nu 2\beta$ Decay



$2\nu 2\beta$ decay ${}^{76}\text{Ge}$: $\tau = 1.5 \times 10^{21} \text{ y}$

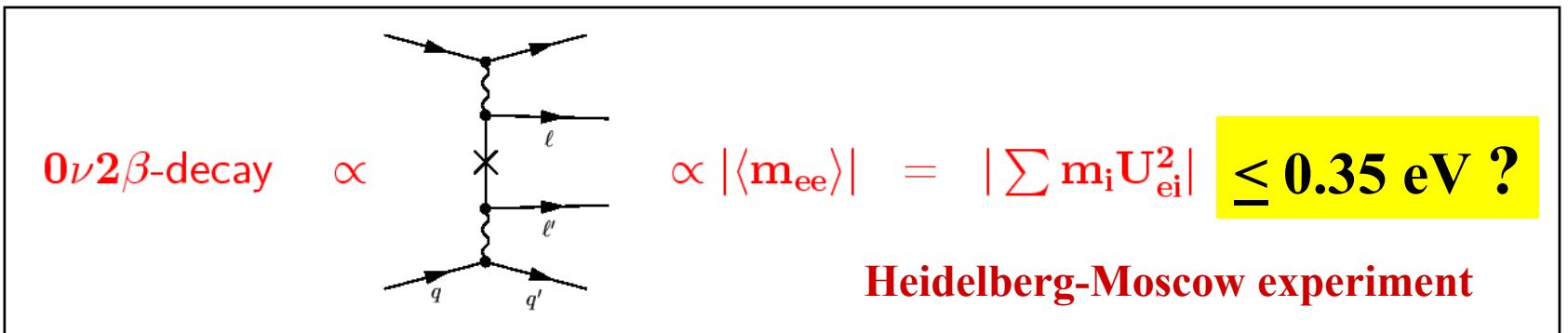


Majorana $\nu \rightarrow 0\nu 2\beta$ decay



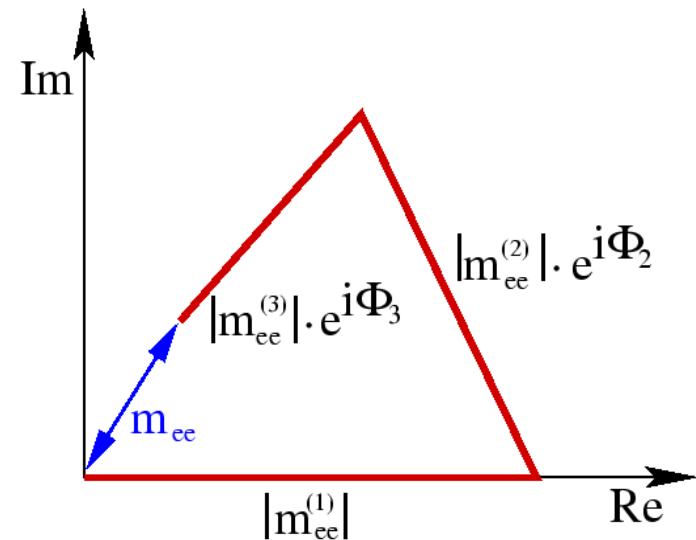
$0\nu 2\beta$
↔
Majorana neutrinos
↔
lepton number violation

Neutrino-less Double β -Decay



$$m_{ee} = |m_{ee}^{(1)}| + |m_{ee}^{(2)}| \cdot e^{i\Phi_2} + |m_{ee}^{(3)}| \cdot e^{i\Phi_3}$$

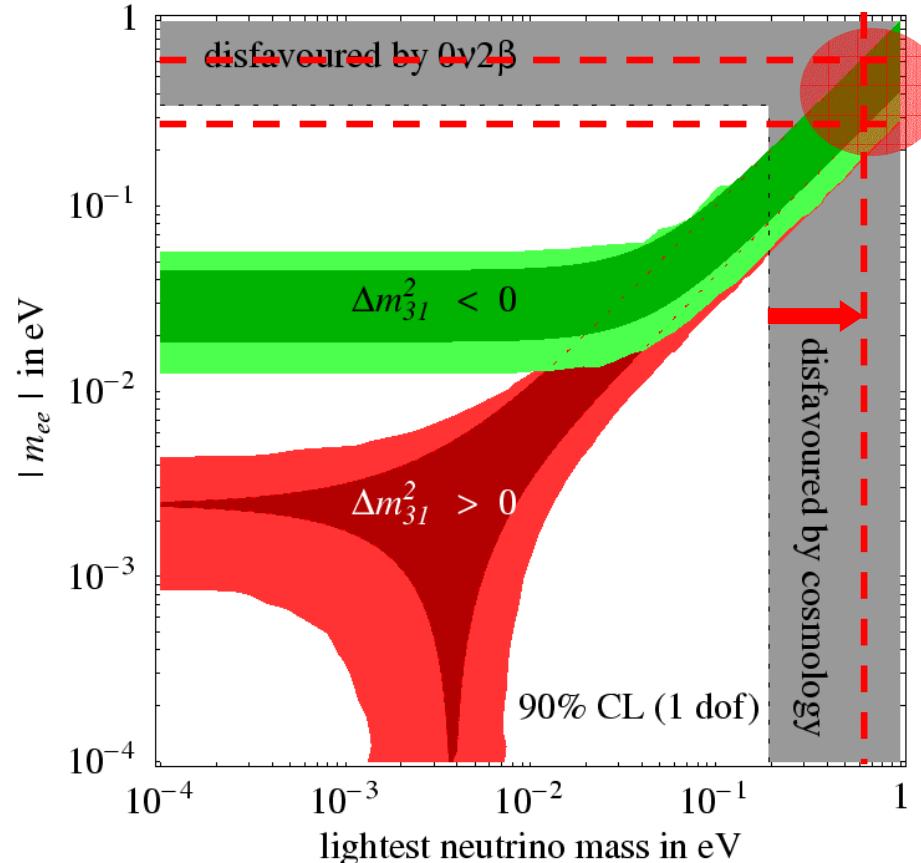
$$\begin{aligned} |m_{ee}^{(1)}| &= |U_{e1}|^2 m_1 \\ |m_{ee}^{(2)}| &= |U_{e2}|^2 \sqrt{m_1^2 + \Delta m_{21}^2} \\ |m_{ee}^{(3)}| &= |U_{e3}|^2 \sqrt{m_1^2 + \Delta m_{31}^2} \end{aligned}$$



solar $\Rightarrow |U_{e1}|^2, |U_{e2}|^2, \Delta m_{21}^2$ atmosph. $\Rightarrow |\Delta m_{31}^2|$ CHOOZ $\Rightarrow |U_{e3}|^2 < 0.05$

→ free parameters: m_1 , sign(Δm_{31}^2), CP-phases Φ_1, Φ_2

- sign of Δm_{31}^2
- partial compensations:
variation of Φ_i , $|U_{e3}|^2$, ...
- model dependent:
LR, SUSY R-parity violation, ...
- nuclear matrix elements ...
- $|m_{ee}| \Rightarrow$ one CP-phase comb.
 \Rightarrow 2nd phase?



→better experiments are important ↔ L-violating effects

sensitivity aim: $|\langle m \rangle| \geq 10^{-2}$ eV \leftrightarrow inverted hierarchy & $\sqrt{|\Delta m_{31}^2|} \approx 0.05$ eV

New Experiments

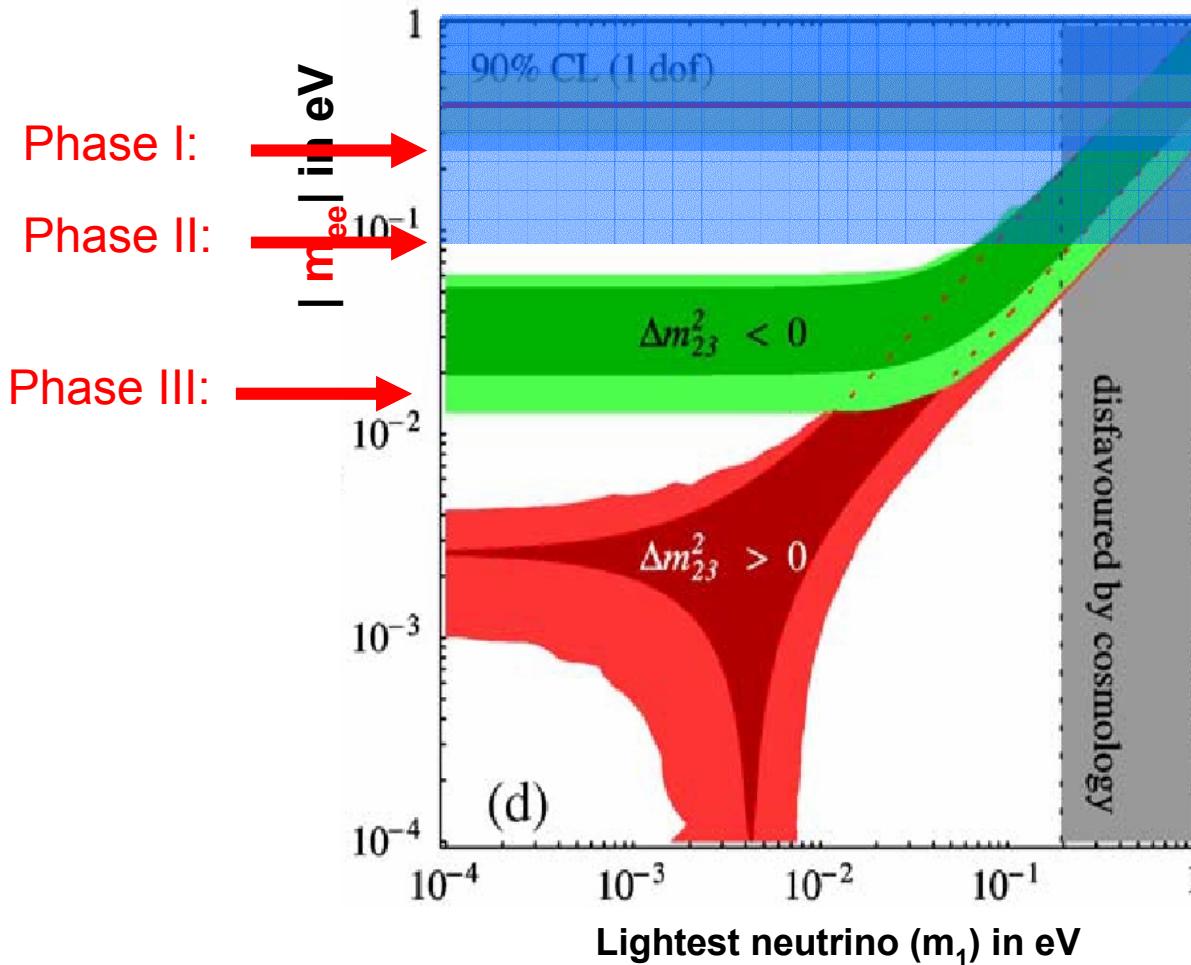
Approved:

- CUORE (Te-130)
- GERDA (Ge-76)

Planned:

- Majorana
- EXO
- MOON
- Super-NEMO
- ...

Phases and Physics reach of GERDA



F.Feruglio, A. Strumia, F. Vissani, NPB 659

ME from Faessler → P-I: 0.31 eV, P-II: 0.12 eV; P-III: 0.02 eV

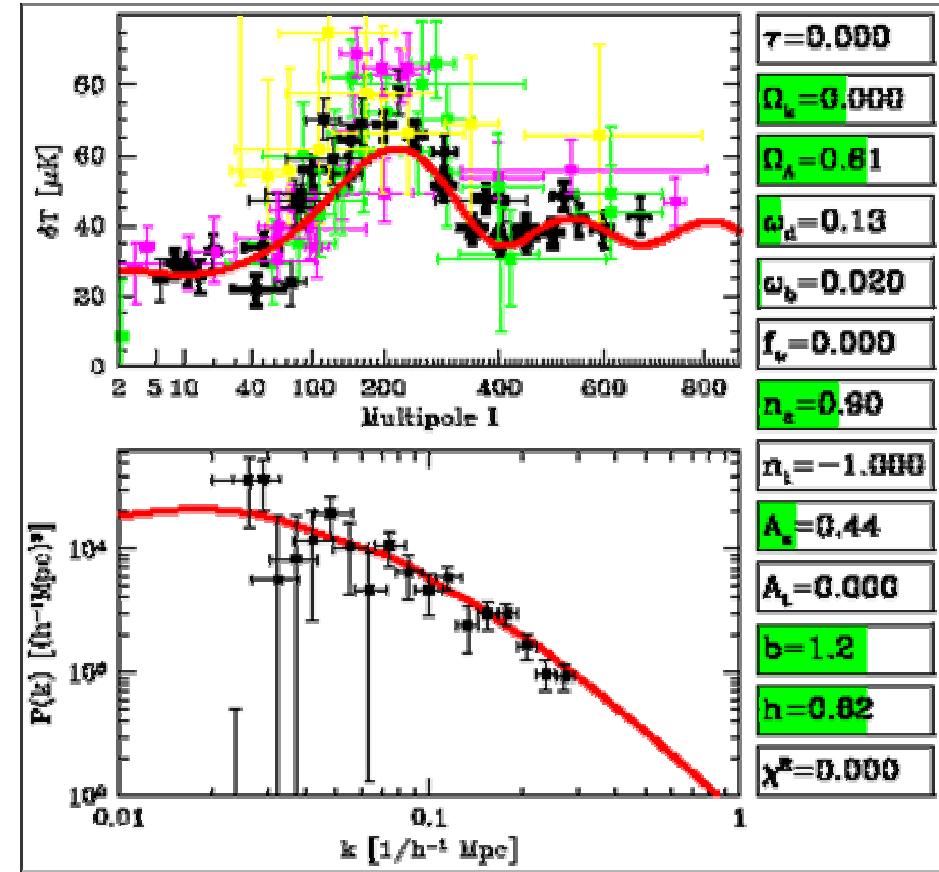
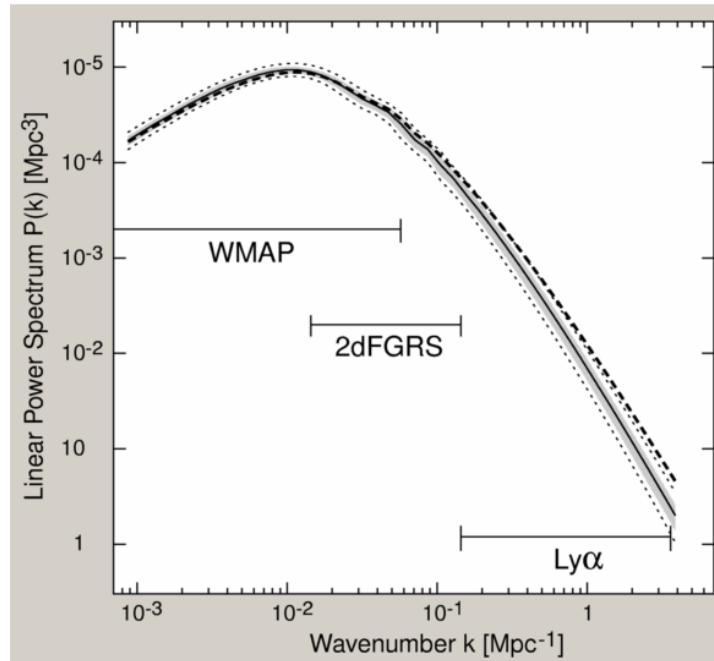
4 Ways to measure Masses & Mixings

4 different methods:

- kinematical
- lepton number violation \leftrightarrow Majorana
- astrophysics & cosmology \rightarrow see next section
- oscillations

Cosmology and Neutrino Mass

- ν 's are hot dark matter → smears structure formation on small scales



• WMAP+2dFGRS + Ly α

→ mass bound: $\sum m_\nu < 0.7 \dots 1.2 \text{ eV}$

• 3 degenerate neutrinos

→ $m_\nu < 0.4 \text{ eV}$ future improvements: ~factor 5...

• comparison with KATRIN, 0ν2β, LSND

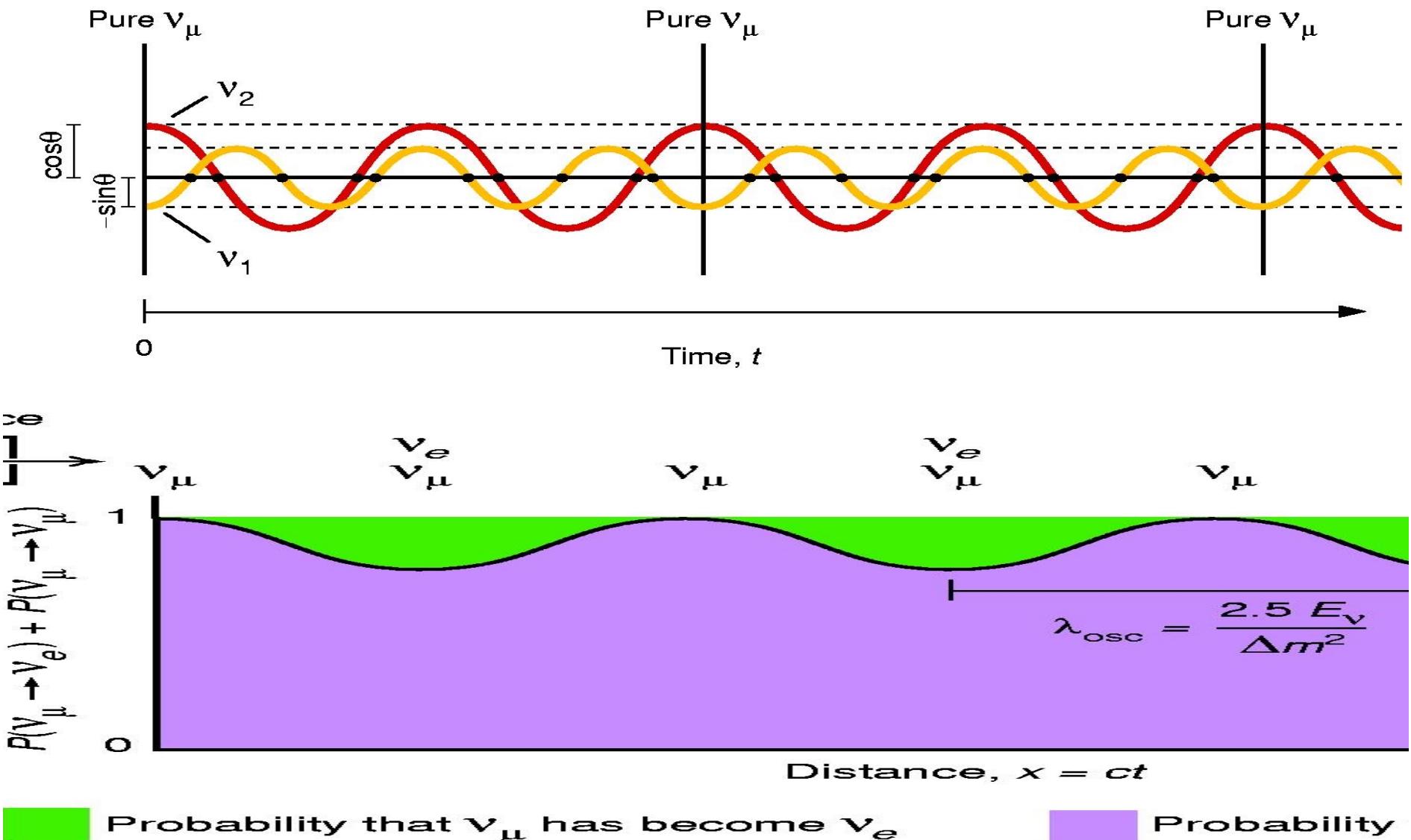
$$f_\nu = \Omega_\nu / \Omega_{\text{matter}}$$

4 Ways to measure Masses & Mixings

4 different methods:

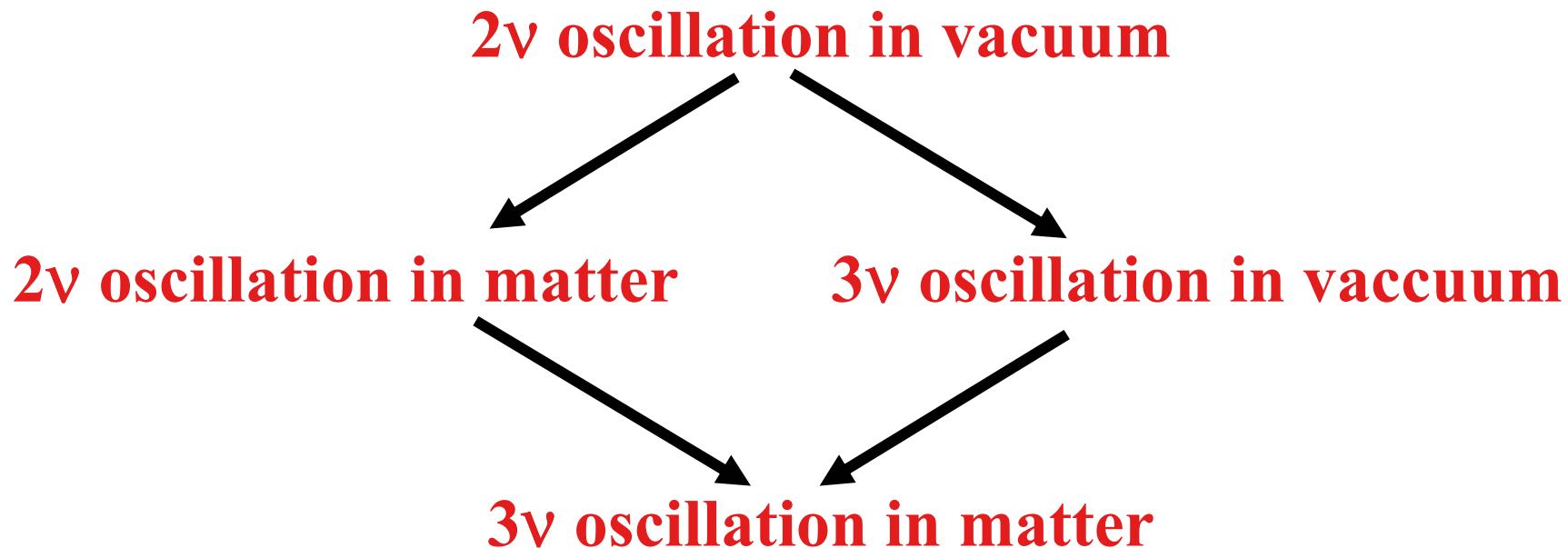
- kinematical
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Neutrino Oscillations



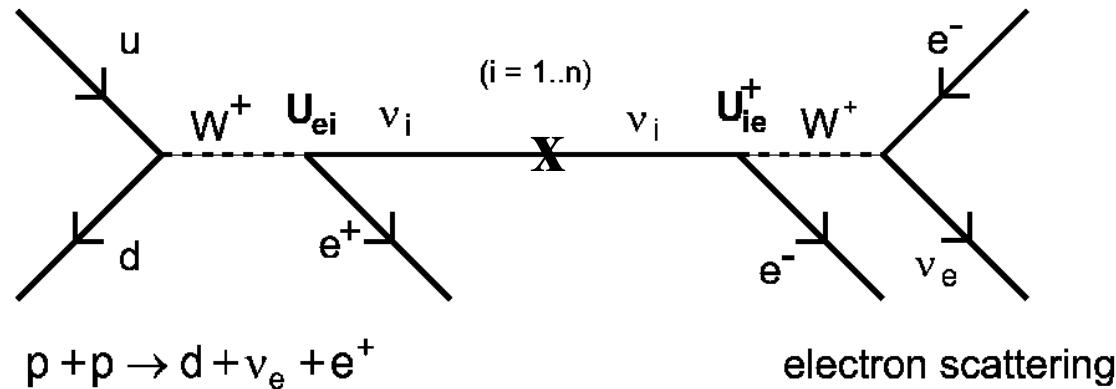
Different Levels of Complexity

Oscillations can involve **2 or 3 neutrinos in vacuum or matter**
→ 4 possibilities:



Neutrino Oscillations for N=2

Production
Propagation
Detection



2 Flavours ν_e, ν_μ :

$$\begin{pmatrix} \nu_e \\ \nu_\mu \end{pmatrix} = \begin{pmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{pmatrix} \cdot \begin{pmatrix} \nu_1 \\ \nu_2 \end{pmatrix}$$

flavour states

mixing matrix

· mass eigenstates

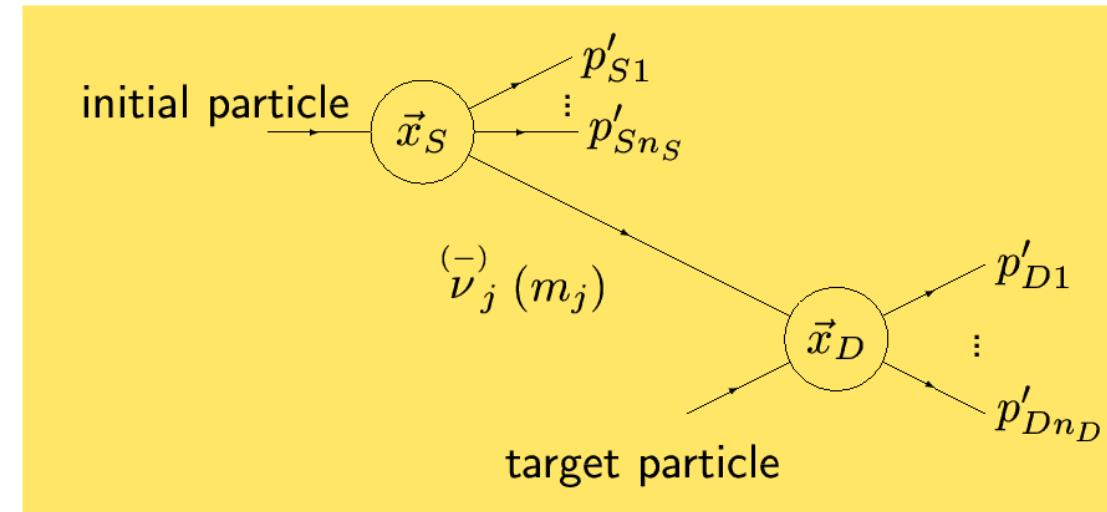
- Production as flavour eigenstate from W-exchange
- Detection via W-exchange \equiv projection on flavour state
- Propagation as mass eigenstate \Rightarrow use mixing matrix at vertex
- Is a simple QM treatment justified?

Neutrino Oscillations in QFT

QFT description of a neutrino produced in a decay at rest:

- localized source and detector
- $L = |\vec{x}_D - \vec{x}_S|$
- initial particle at rest
- target particle at rest

... DIF similar



Transition probability from Feynman diagram:

$$\left\langle P_{(\bar{\nu})_\alpha \rightarrow (\bar{\nu})_\beta} \right\rangle_{\mathcal{P}} \propto \int dP_S \int_{\mathcal{P}} \frac{d^3 p_{D1}}{2E_{D1}} \dots \frac{d^3 p_{Dn_D}}{2E_{Dn_D}} \left| \mathcal{A}_{(\bar{\nu})_\alpha \rightarrow (\bar{\nu})_\beta} \right|^2$$

Neutrino propagator in ultra-relativistic limit $\rightarrow e^{ipx}$

QFT derivation: Avoids confusion (factors of two..., etc.)

Kinematics: Equal Energy or equal Momenta?

- Consider e.g. pion decay at rest: $\pi^+ \rightarrow \mu^+ + \nu_\mu$
- Neutrino energy and momentum determined by energy-momentum conservation

$$p_k^2 = \frac{m_\pi^2}{4} \left(1 - \frac{m_\mu^2}{m_\pi^2}\right)^2 - \frac{m_k^2}{2} \left(1 + \frac{m_\mu^2}{m_\pi^2}\right) + \frac{m_k^4}{4m_\pi^2}$$

$$E_k^2 = \frac{m_\pi^2}{4} \left(1 - \frac{m_\mu^2}{m_\pi^2}\right)^2 + \frac{m_k^2}{2} \left(1 - \frac{m_\mu^2}{m_\pi^2}\right) + \frac{m_k^4}{4m_\pi^2}$$

- For $E \gg m$:
$$p_k \simeq E - \xi \frac{m_k^2}{2E}, \quad E_k \simeq E + (1 - \xi) \frac{m_k^2}{2E}$$

with $E = \frac{m_\pi}{2} \left(1 - \frac{m_\mu^2}{m_\pi^2}\right) \simeq 30 \text{ MeV}$, $\xi = \frac{1}{2} \left(1 + \frac{m_\mu^2}{m_\pi^2}\right) \simeq 0.8$

\Rightarrow neither equal energy nor equal momentum!

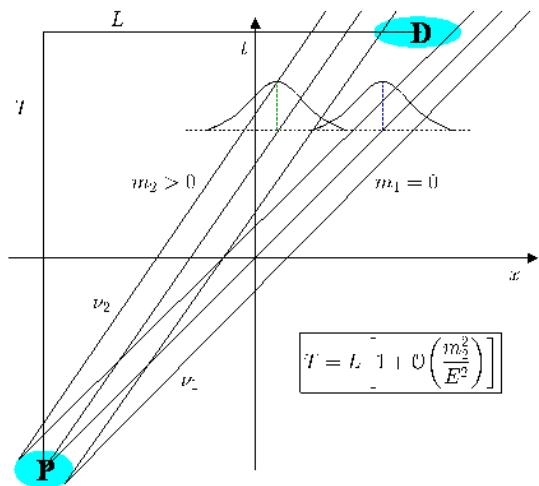
$$e^{ipx} \Rightarrow \boxed{p_\mu \cdot x^\mu = p_k L - E_k T = -\frac{m_k^2 L}{2E}} \text{ for } L = T$$

\Rightarrow ξ drops out of the oscillation formulae \Leftrightarrow naive treatment correct

- Shown for π -decay, but valid in general (DIF, N-body, ..., different ξ)

Localized Source and Detector:

- Feynman rules for particles of given momentum (\simeq on-shell)
 \Rightarrow this corresponds to an infinitely extended (non-localized) plane wave
- Localized source (wave packet) and detector in space-time $(\Delta x_S, \Delta t_S)$, $(\Delta x_D, \Delta t_D)$:
 \Rightarrow Source: Fourier superposition of momenta with $\sigma_S^2 \simeq \min(\Delta x_S^2, \Delta t_S^2)$
 \Rightarrow Detector: projection on a superposition of momenta with $\sigma_D^2 \simeq \min(\Delta x_D^2, \Delta t_D^2)$
- Different masses and momenta \Rightarrow dispersion \Rightarrow loss of coherence



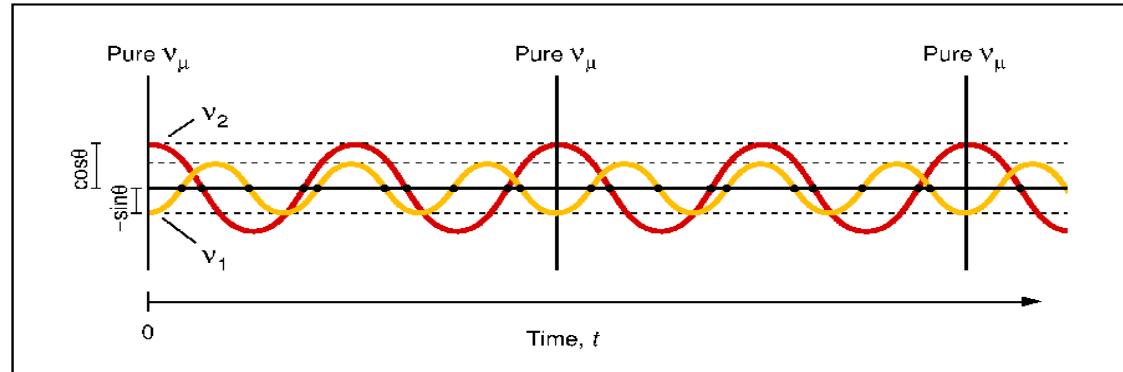
- separation of wave packets:
- $$\Delta x_{ij} = \frac{|\Delta m_{ij}|L}{2E^2}$$
- coherence condition: $\Delta x_{ij} < \sigma := \sqrt{\sigma_S^2 + \sigma_D^2}$
 - coherence length: $L < \frac{2E^2\sigma}{|\Delta m_{ij}|}$

- Oscillations from QFT $\Rightarrow P_{\nu_\alpha \rightarrow \nu_\beta}(L, T) = \left| \sum_k U_{\alpha k}^* e^{ip_k L - iE_k T} U_{\beta k} \right|^2$
- Very interesting QM effects (σ , decay)

Two Neutrino Oscillations

2 Neutrinos: ν_e, ν_μ

$$\begin{aligned} |\nu_e(0)\rangle &= \cos\theta |\nu_1\rangle + \sin\theta |\nu_2\rangle \\ |\nu_\mu(0)\rangle &= -\sin\theta |\nu_1\rangle + \cos\theta |\nu_2\rangle \end{aligned}$$



$$|\nu_\mu(t)\rangle = -\sin\theta \exp[-\frac{iE_1t}{\hbar}] |\nu_1\rangle + \cos\theta \exp[-\frac{iE_2t}{\hbar}] |\nu_2\rangle$$

$$E_i = \sqrt{p_i^2 + m_i^2} \xrightarrow{p_i=p \gg m_i} \simeq p + \frac{m_i^2}{2p} \simeq p + \frac{m_i^2}{2E}$$

$$L = c \cdot t \quad \Delta m^2 = m_2^2 - m_1^2 \Rightarrow E_2 - E_1 = \frac{\Delta m^2}{2E}$$

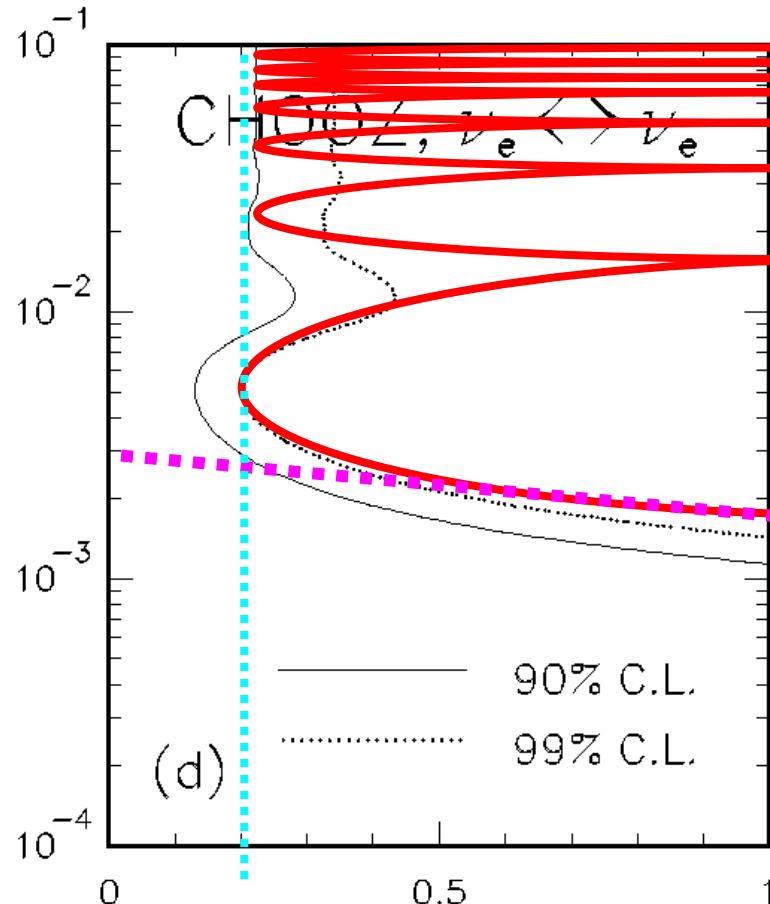
**2ν-transition-
probability:**

$$P(\nu_\mu \rightarrow \nu_e) = |\langle \nu_\mu(t) | \nu_e(0) \rangle|^2 = \sin^2 2\theta \cdot \sin^2 \left(\frac{\Delta m^2 L}{4E} \right)$$

$\nu_e, \nu_\mu, \nu_\tau \rightarrow$ 9 oscillation channels for neutrinos
 $\nu_e, \nu_\mu, \nu_\tau \rightarrow$ 9 channels for anti-neutrinos (assuming 3ν !)

2 Flavour Exclusion Limits

$$P(\nu_\mu \rightarrow \nu_e) = \sin^2(2\theta) \cdot \sin^2\left(\frac{\Delta m^2 L}{4E}\right)$$



small Δm^2 :

$$P \simeq \sin^2(2\theta) \left(\frac{\Delta m^2 L}{4E}\right)^2 < P_{exp.}$$

$$\rightarrow \sin(2\theta) \cdot \Delta m^2 < \frac{\sqrt{P}4E}{L}$$

$\sin(\Delta m^2 L / 4E) = 1$:

$$P \simeq \sin^2(2\theta) < P_{exp.}$$

$$\sin(2\theta) < \sqrt{P}$$

Inbetween: Oscillation + energy smearing

Other Oscillation Phenomena

Possibilities:

- different neutrino flavours
- $B_0 \leftrightarrow \bar{B}_0$ $K_0 \leftrightarrow \bar{K}_0$
- neutrinos \leftrightarrow anti-neutrinos
- neutron \leftrightarrow anti-neutron
- different charged lepton flavours
- within up and down quark flavours

Requirements:

- ✓
- ✓
- ✓ new physics
- ✓ new physics
- ✓ new basis...
- ✓ distance...

Crutial:

- identical quantum numbers
- mixing

$N \geq 2$ with CP Violation and Matter Effects

Precision: $N = 2$ description insufficient \Rightarrow modifications

- $2 \rightarrow 3$ neutrino framework \Rightarrow more parameters & CP effects
- MSW: parameter mapping in matter

Quantum mechanical treatment for N ultra-relativistic neutrinos:

$$P_{\nu_{e_l} \rightarrow \nu_{e_m}}(L/E) = \left| \sum_j U_{mj} U_{lj}^* \exp\left(\frac{-i m_j^2 L}{2E}\right) \right|^2$$

- masses m_j associated to mass eigenfields ν_j
- flavour eigenstates $\nu_{e_l}, \nu_{e_m}, \dots$
- neutrino energy E
- source and detector distance L
- unitary mixing matrix U

General 3x3 neutrino mixing matrix:

has (up to) 3 angles + 1 Dirac-phase **+2 Majorana-phases**: θ_{12} , θ_{23} , θ_{13} , δ , Φ_1 , Φ_2

$$U_{MNS} = \boxed{U \cdot \text{diag}(\exp[i\Phi_1], \exp[i\Phi_2], 1)}$$

$$U = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix}$$

- Only U enters in neutrino oscillations:

$$J_{ij}^{e_l e_m} := U_{li} U_{lj}^* U_{mi}^* U_{mj}$$

- All oscillation frequencies show up:

$$\Delta_{ij} := \frac{\Delta m_{ij}^2 L}{4E} = \frac{(m_i^2 - m_j^2)L}{4E}$$

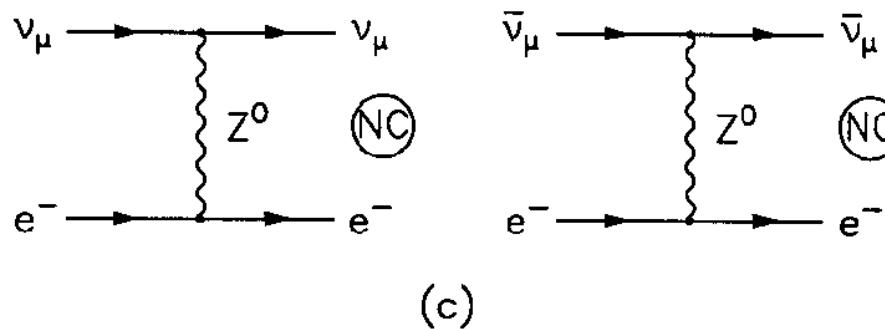
$$P(\nu_{e_l} \rightarrow \nu_{e_m}) = \underbrace{\delta_{lm} - 4 \sum_{i>j} \text{Re} J_{ij}^{e_l e_m} \sin^2 \Delta_{ij}}_{P_{CP}} - \underbrace{2 \sum_{i>j} \text{Im} J_{ij}^{e_l e_m} \sin 2\Delta_{ij}}_{P_{CP}}$$

\Rightarrow **Leptonic CP violation, genuine 3 flavour and matter effects**

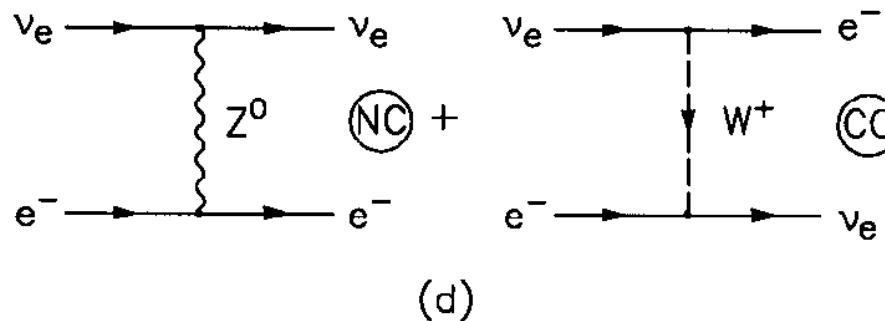
→ lengthy 3 flavour oscillation formulae

Matter Effects and MSW Resonance

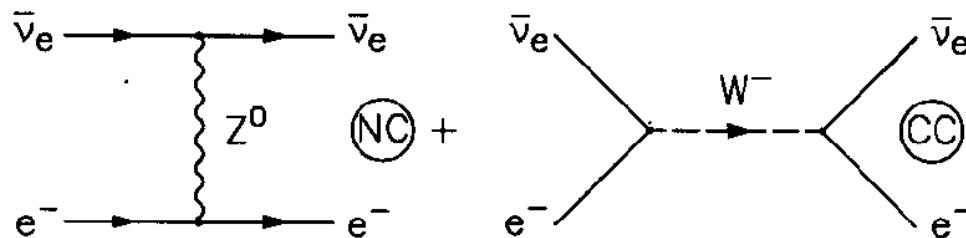
Mikheyev-Smirnov-Wolfenstein: coherent forward scattering



(c)



(d)



\mathcal{L}_{NC} = flavour universal

$\mathcal{L}_{CC} = \sqrt{2}G_F n_e \Leftrightarrow$ only ν_e

MSW-resonance energy (Δm_{31}^2)

Earth: $E_{res} \simeq 10$ GeV

for beams
dominated by average density

$$\rho = \rho_{\text{average}} + \delta\rho$$

$$\mathcal{L}_{CC} = -\frac{G_F}{\sqrt{2}} [\bar{e}\gamma_\mu(1-\gamma_5)\nu_e] [\bar{\nu}_e\gamma^\mu(1-\gamma_5)e]$$

Fierz

$$\textbf{Transf.} = -\frac{G_F}{\sqrt{2}} [\bar{e}\gamma_\mu(1-\gamma_5)e] [\bar{\nu}_e\gamma^\mu(1-\gamma_5)\nu_e]$$

$$H_{eff}(\nu) = \langle \mathcal{L}_{eff} \rangle_e = \bar{\nu}_e \ V_{eff} \ \nu_e$$

$$\langle \bar{e}\gamma_o e \rangle = \langle e^+ e \rangle = N_e$$

$$\langle \bar{e}\gamma_i e \rangle; \langle \bar{e}\gamma_o\gamma_5 e \rangle; \langle \bar{e}\gamma_o\gamma_5 e \rangle \simeq 0$$

$$V_{eff}(\nu_e) = \sqrt{2}G_F N_e$$

Hamiltonian for 3 Neutrino Oscillations in Flavour Basis:

$$\mathbf{H} = H_0 + \delta\mathbf{H}_{CC} + \delta\mathbf{H}_{NC} = \frac{1}{2E} \mathbf{U} \begin{pmatrix} m_1^2 & 0 & 0 \\ 0 & m_2^2 & 0 \\ 0 & 0 & m_3^2 \end{pmatrix} \mathbf{U}^T + \frac{1}{2E} \begin{pmatrix} \mathbf{A} + \mathbf{A}' & 0 & 0 \\ 0 & \mathbf{A}' & 0 \\ 0 & 0 & \mathbf{A}' \end{pmatrix}$$

- $\mathbf{A} = \pm \frac{2\sqrt{2}G_F Y \rho E}{m_n} = 2V \cdot E$ $\nu \oplus \text{matter}$ and $\bar{\nu} \oplus \text{anti-matter} \Rightarrow$ “+”
- $Y = e^-/\text{nucleon}$ $\rho = \text{matter density}$ $m_n = \text{nucleon mass}$
- Overall phases drop out: $m_i \rightarrow m_i - m_1 \Rightarrow m_1$ and A' can be eliminated

$$\mathbf{H}' = \frac{1}{2E} \mathbf{U} \begin{pmatrix} 0 & 0 & 0 \\ 0 & \Delta m_{21}^2 & 0 \\ 0 & 0 & \Delta m_{31}^2 \end{pmatrix} \mathbf{U}^T + \frac{1}{2E} \begin{pmatrix} \mathbf{A} & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

Note: A' can not be eliminated if sterile neutrinos are present

- In good approximation $\Delta m_{12}^2 \simeq 0$
- \mathbf{U} can be written as a sequence of rotations: $\mathbf{U} = R_{23}R_{13}R_{12}$

$$\begin{aligned}
\mathbf{H}'' &= \frac{1}{2E} \mathbf{R}_{23} \mathbf{R}_{13} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & \Delta m_{31}^2 \end{pmatrix} \mathbf{R}_{13}^{-1} \mathbf{R}_{23}^{-1} + \frac{1}{2E} \mathbf{R}_{23} \begin{pmatrix} \mathbf{A} & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \mathbf{R}_{23}^{-1} \\
&= \frac{1}{2E} \mathbf{R}_{23} \left[\mathbf{R}_{13} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & \Delta m_{31}^2 \end{pmatrix} \mathbf{R}_{13}^{-1} + \begin{pmatrix} \mathbf{A} & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \right] \mathbf{R}_{23}^{-1} \\
&= \frac{1}{2E} \mathbf{R}_{23} \left[\begin{pmatrix} \bullet & 0 & \bullet \\ 0 & 0 & 0 \\ \bullet & 0 & \bullet \end{pmatrix} + \begin{pmatrix} \mathbf{A} & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \right] \mathbf{R}_{23}^{-1} \\
&= \frac{1}{2E} \mathbf{R}_{23} \left[\mathbf{R}'_{13} \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & \Delta(m_{31}^2)' \end{pmatrix} (\mathbf{R}'_{13})^{-1} \right] \mathbf{R}_{23}^{-1}
\end{aligned}$$

\Rightarrow re-insert $R_{12} \Rightarrow U' \Rightarrow$ parameter mapping in 1-3 subspace

- **Different mappings for neutrinos and antineutrinos**
- 1-3 sub-space mapping like in 2 neutrino case

- Relevant quantity

$$C_{\pm}^2 = \left(\frac{A}{\Delta m_{31}^2} - \cos 2\theta_{13} \right)^2 + \sin^2 2\theta_{13}$$

- MSW resonance condition for $\theta_{13} \simeq 0$: $\Delta m_{31}^2 = A = 2VE = \pm \frac{2\sqrt{2}G_F Y \rho E}{m_n}$
- Effective parameters in matter:

$$\begin{aligned}\sin^2 2\theta'_{13} &= \frac{\sin^2 2\theta_{13}}{C_{\pm}^2} \\ \Delta m_{31,m}^2 &= \Delta m^2 C_{\pm} \\ \Delta m_{32,m}^2 &= \frac{\Delta m^2 (C_{\pm} + 1) + A}{2} \\ \Delta m_{21,m}^2 &= \frac{\Delta m^2 (C_{\pm} - 1) - A}{2}\end{aligned}$$

- Corrections due to
 - $\Delta m_{12}^2 \neq 0$
 - non-constant matter profiles \Rightarrow solve Schrödinger equation

Analytic Approximations

- $\Delta = \Delta m_{31}^2 L / 4E$
- qualitative understanding \Rightarrow expand in $\alpha = \Delta m_{21}^2 / \Delta m_{31}^2$ and $\sin^2 2\theta_{13}$
- matter effects $\hat{A} = A / \Delta m_{31}^2 = 2VE / \Delta m_{31}^2$; $V = \sqrt{2}G_F n_e$

$$P(\nu_\mu \rightarrow \nu_\mu) \approx 1 - \cos^2 \theta_{13} \sin^2 2\theta_{23} \sin^2 \Delta + 2 \alpha \cos^2 \theta_{13} \cos^2 \theta_{12} \sin^2 2\theta_{23} \Delta \cos \Delta$$

$$\begin{aligned} P(\nu_e \rightarrow \nu_\mu) &\approx \sin^2 2\theta_{13} \sin^2 \theta_{23} \frac{\sin^2((1-\hat{A})\Delta)}{(1-\hat{A})^2} \\ &\pm \sin \delta_{\text{CP}} \alpha \sin 2\theta_{12} \cos \theta_{13} \sin 2\theta_{13} \sin 2\theta_{23} \sin(\Delta) \frac{\sin(\hat{A}\Delta) \sin((1-\hat{A})\Delta)}{\hat{A}(1-\hat{A})} \\ &+ \cos \delta_{\text{CP}} \alpha \sin 2\theta_{12} \cos \theta_{13} \sin 2\theta_{13} \sin 2\theta_{23} \cos(\Delta) \frac{\sin(\hat{A}\Delta) \sin((1-\hat{A})\Delta)}{\hat{A}(1-\hat{A})} \\ &+ \alpha^2 \sin^2 2\theta_{12} \cos^2 \theta_{23} \frac{\sin^2(\hat{A}\Delta)}{\hat{A}^2} \end{aligned}$$

→ analytic discussion / full simulations

→ degeneracies, correlations → $(\sin^2 2\theta_{13})_{\text{eff}}$

Cervera et al.

Freund, Huber, ML

Akhmedov, Johansson, ML, Ohlsson, Schwetz

Degeneracies, Correlations, ...

Fixed L/E → probabilities invariant under transformations:

- $\theta_{23} \rightarrow \pi/2 - \theta_{23}$ Fogli, Lisi
- $P(v_e \rightarrow v_\mu)$ not really invariant → compensation by small parameter off-sets
- $\Delta m^2 \rightarrow -\Delta m^2$ compensated by offset in δ Minakata, Nunokawa
- $P(v_e \rightarrow v_\mu) = \text{const.} \rightarrow \delta - \theta_{13}$ manifolds Koike, Ota, Sato & Burguet-Castell et al.
- → 8-fold degeneracy Barger, Marfatia, Whisnant

- parameter extraction suffers from correlations & degeneracies
- how to break degeneracies & correlations?

Baseline & MSW Matter Effect

Beams in Earth Matter:

⇒ electron density profile
as function of radius **Stacy**
density errors **Geller & Hara**

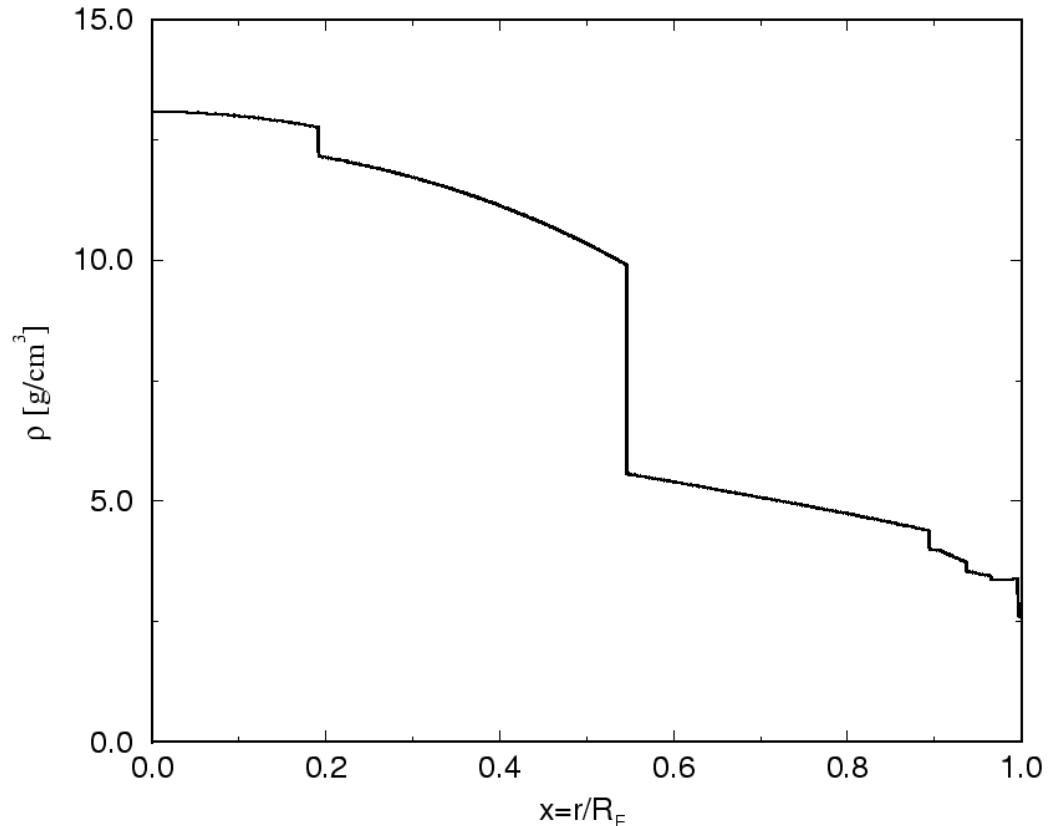
Large $L \Rightarrow$ steep angles

$$L = 2800 \text{ km} \Leftrightarrow 13^\circ$$

$$L = 7300 \text{ km} \Leftrightarrow 35^\circ$$

$$L = 12750 \text{ km} \Leftrightarrow 90^\circ$$

$L \leq \mathcal{O}(10\,000 \text{ km}) \Leftrightarrow$ **mantle**



- $E_{\text{resonance}} \simeq 10 - 15 \text{ GeV}$, matter effects grow with distance L
- Average density profile uncertainties decrease with $L \Rightarrow \simeq 5\%$ error

The Magic Baseline

$$\begin{aligned}
 P(\nu_e \rightarrow \nu_\mu) &\approx \sin^2 2\theta_{13} \sin^2 \theta_{23} \frac{\sin^2((1-\hat{A})\Delta)}{(1-\hat{A})^2} \\
 &\pm \sin \delta_{\text{CP}} \alpha \sin 2\theta_{12} \cos \theta_{13} \sin 2\theta_{13} \sin 2\theta_{23} \sin(\Delta) \frac{\sin(\hat{A}\Delta) \sin((1-\hat{A})\Delta)}{\hat{A}(1-\hat{A})} \\
 &+ \cos \delta_{\text{CP}} \alpha \sin 2\theta_{12} \cos \theta_{13} \sin 2\theta_{13} \sin 2\theta_{23} \cos(\Delta) \frac{\sin(\hat{A}\Delta) \sin((1-\hat{A})\Delta)}{\hat{A}(1-\hat{A})} \\
 &+ \alpha^2 \sin^2 2\theta_{12} \cos^2 \theta_{23} \frac{\sin^2(\hat{A}\Delta)}{\hat{A}^2}
 \end{aligned}$$

- All terms besides the first vanish for $\sin(\hat{A}\Delta) = 0$
 - Condition for uncorrelated sensitivity to θ_{13} $\boxed{\hat{A}\Delta = \pi}$
- \Rightarrow inserting $\hat{A} = A/\Delta m_{31}^2$, $A = 2VE$, $\Delta = \Delta m_{31}^2 L/4E$ one finds

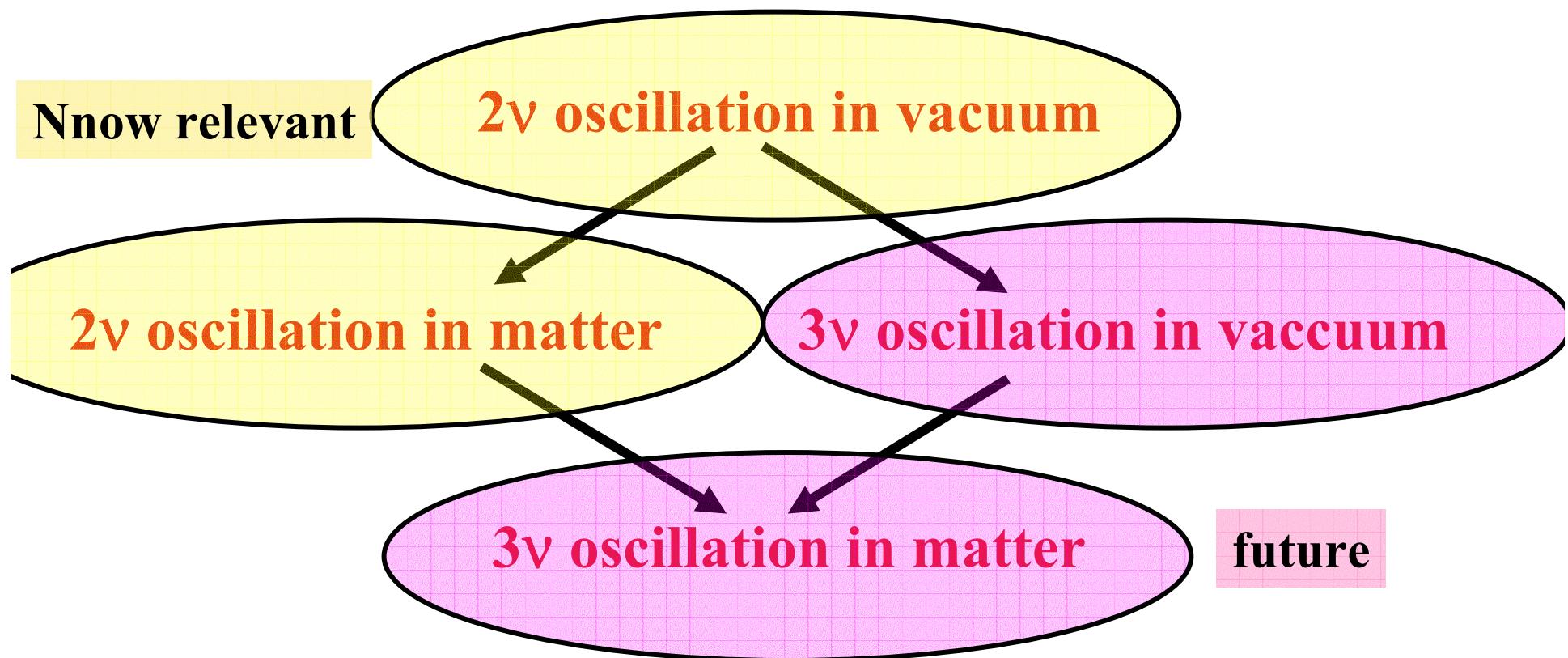
$$L_{\text{magic}} = \frac{2\pi}{\sqrt{2}G_F n_e} = 7630 \text{ km} \cdot \frac{\rho}{4.3g/cm^3}$$

Huber, Winter

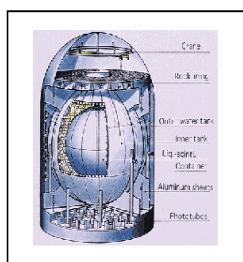
- Note that this is not the MSW resonance condition

Different Levels of Complexity

Oscillations can involve **2 or 3 neutrinos** in **vacuum or matter**
→ 4 possibilities:



Neutrino Oscillation Results: Overview



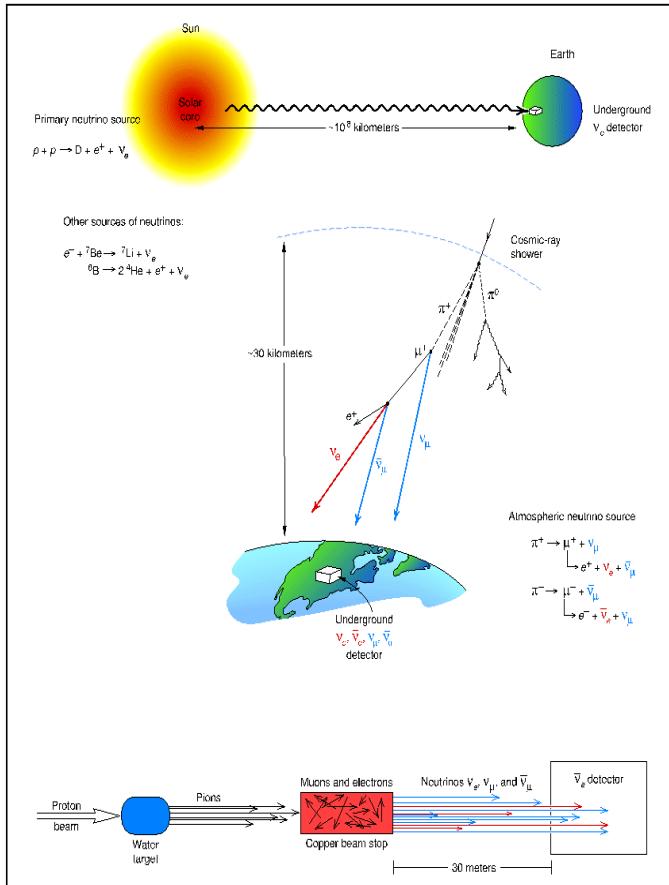
KamLAND
CHOOZ

\Leftrightarrow

atmospheric
+
K2K

neutrino
oscillation
signals

LSND?



→ approximately two 2x2 oscillations

$$\Delta m_{12}^2 = 8.2 \pm 0.3 \times 10^{-5} \text{ eV}^2$$

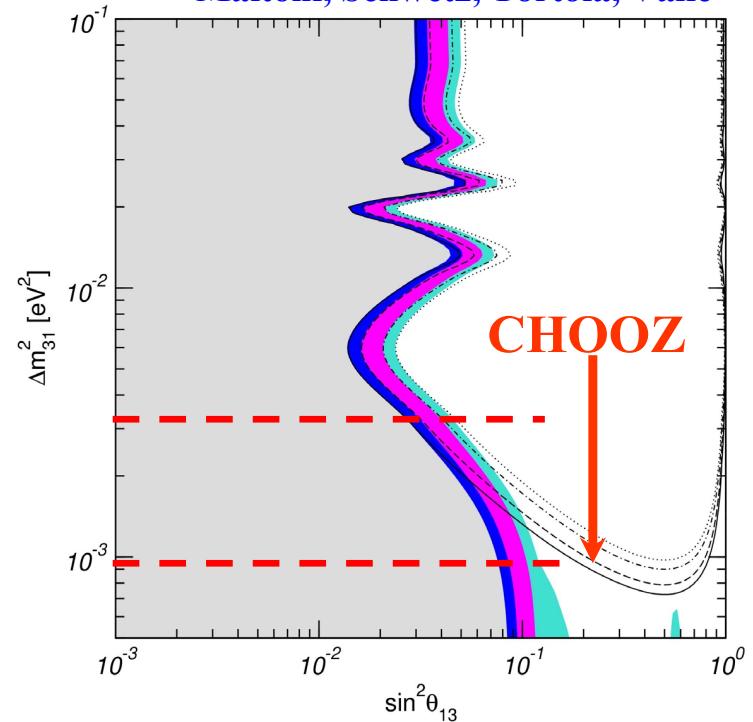
$$\tan^2 \theta_{12} = 0.39 \pm 0.05$$

$$\Delta m_{23}^2 = 2.2 \pm 0.6 \times 10^{-3} \text{ eV}^2$$

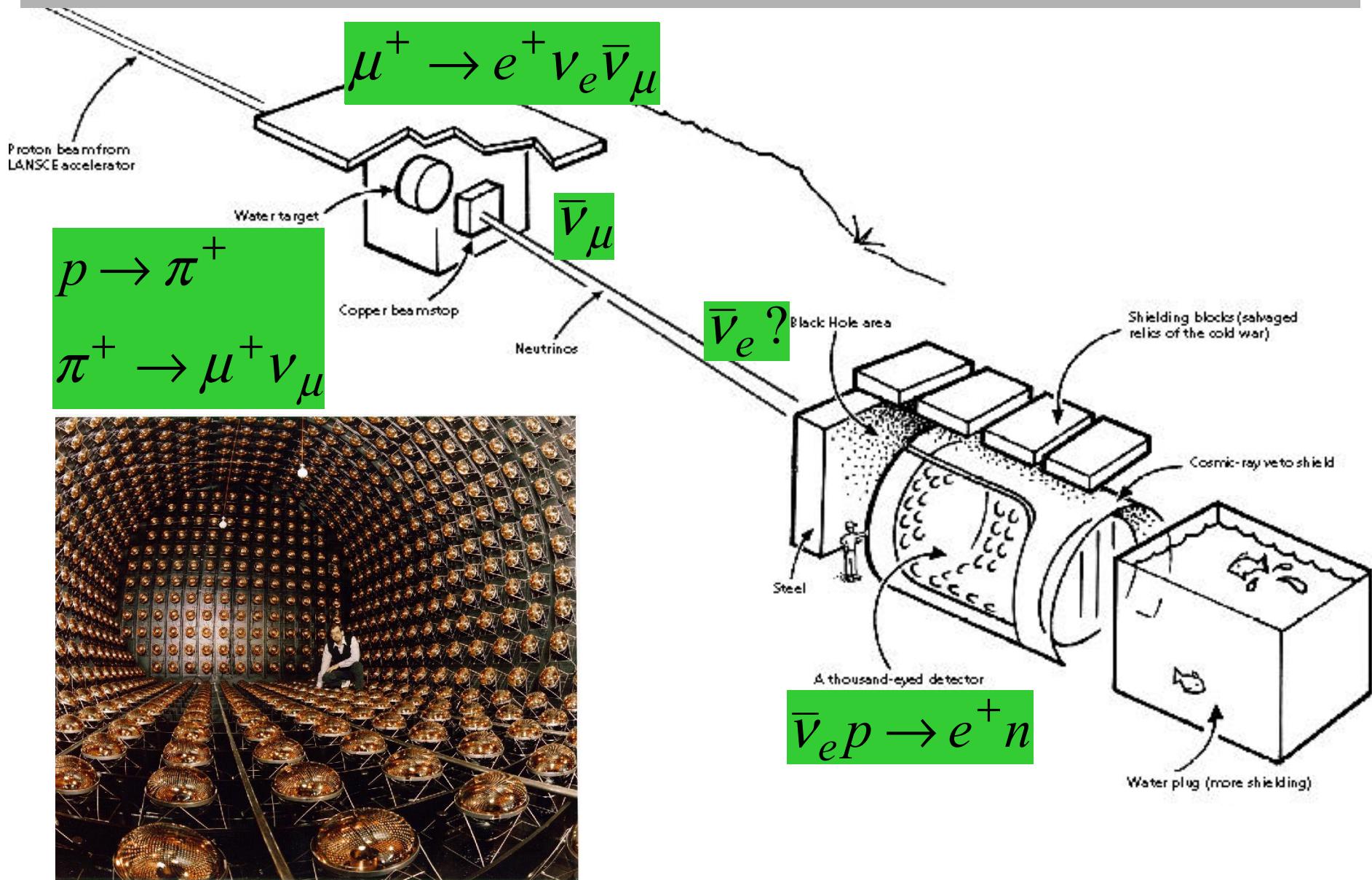
$$\tan^2 \theta_{23} = 1.0 \pm 0.3$$

$\sin^2 \theta_{13} \leq 0.041$ @ 3σ

Maltoni, Schwetz, Tortola, Valle

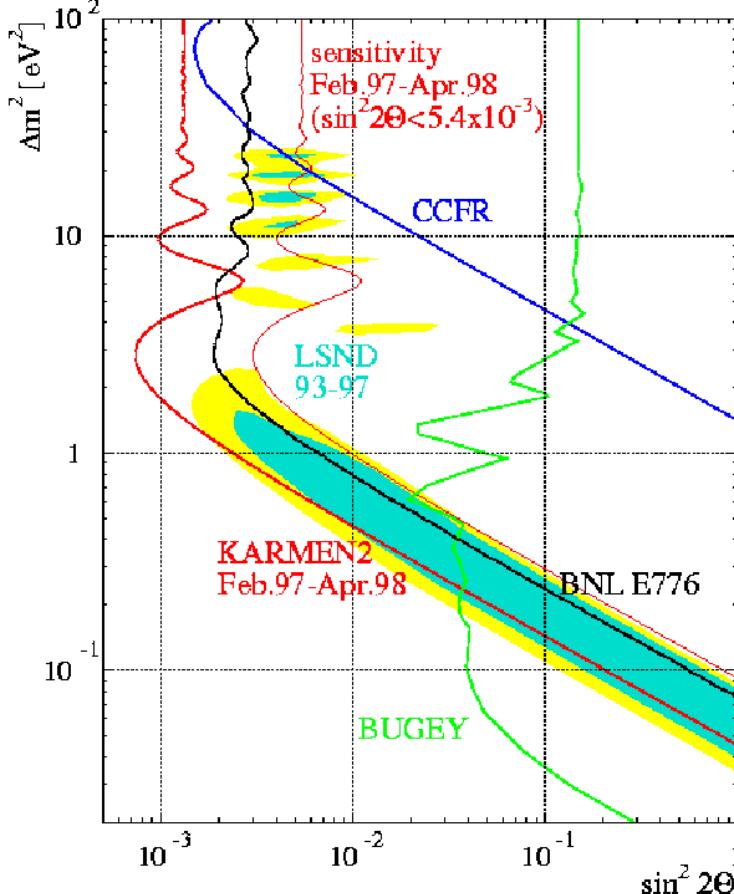


LSND

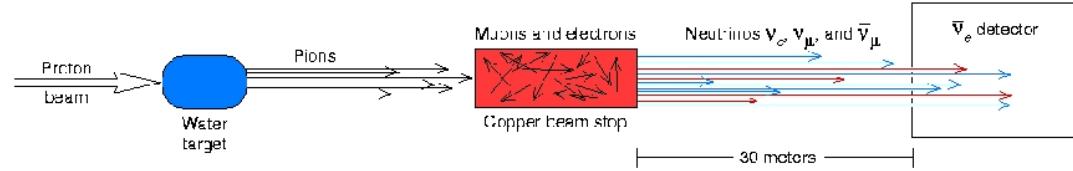


LSND claims Oscillations

KARMEN Oscillation Limit (Unified Approach)
 $\sin^2 2\Theta < 1.3 \times 10^{-3}$ (90% CL.)



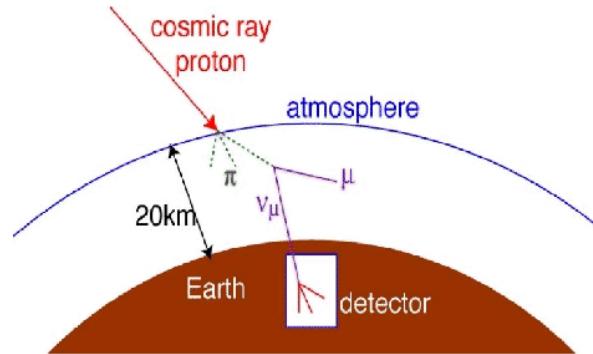
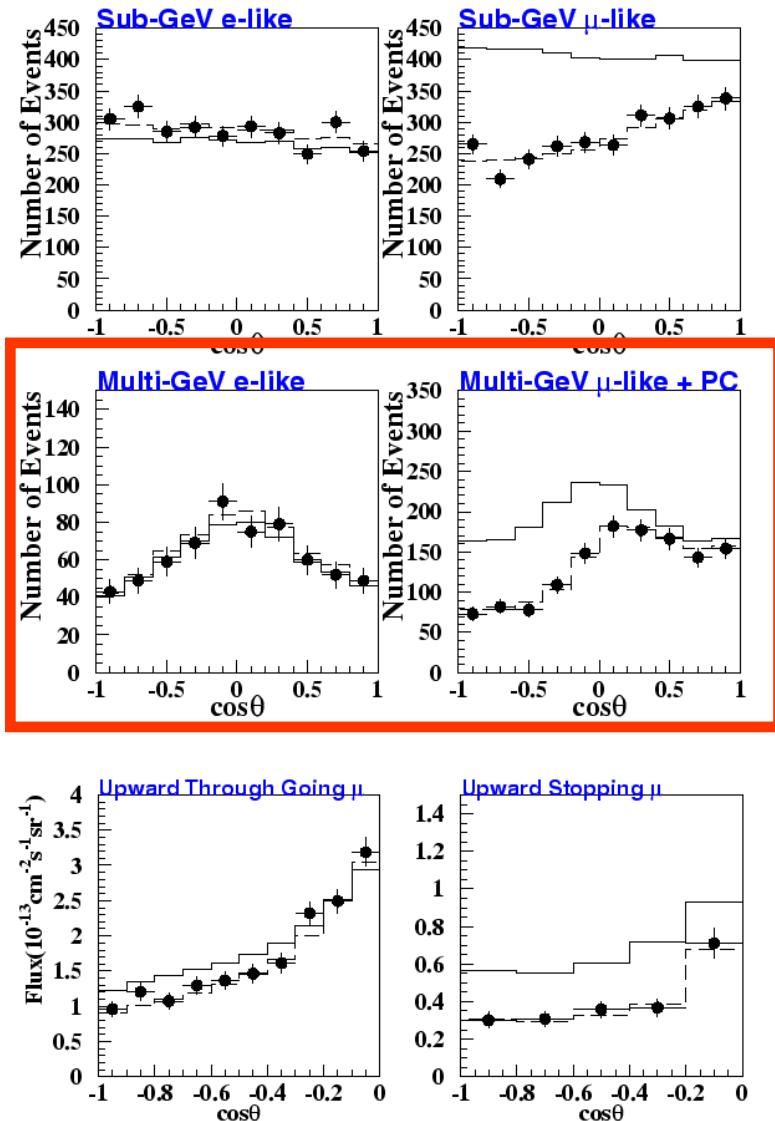
unified approach: R.D. Cousins and G.J. Feldman
 Phys. Rev. D57 (1998) 3873



- LSND claims $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ and $\nu_\mu \rightarrow \nu_e$ osc.
 - 1) Third $\Delta m^2 \simeq 1 \text{ eV} \Rightarrow$ 4 light neutrinos?
 \Leftrightarrow Z line-shape
 - 2) CPT violation = different param. for ν and $\bar{\nu}$
 \Leftrightarrow local QFT!?
- Partly ruled out by KARMEN
- Partly ruled out by \rightarrow cosmology
- Will be tested soon by \Rightarrow MiniBooNE

3+1 \rightarrow tension in global fits
 3+2 OK, 3+N $\leftarrow \rightarrow$ cosmology

Atmospheric Oscillations @SuperK



Oscillation fit (including K2K):

$$|\Delta m_{31}^2| = (2.0 \pm 1.0) \cdot 10^{-3} \text{ eV}^2$$

$$\sin^2 2\theta_{23} \geq 0.85$$

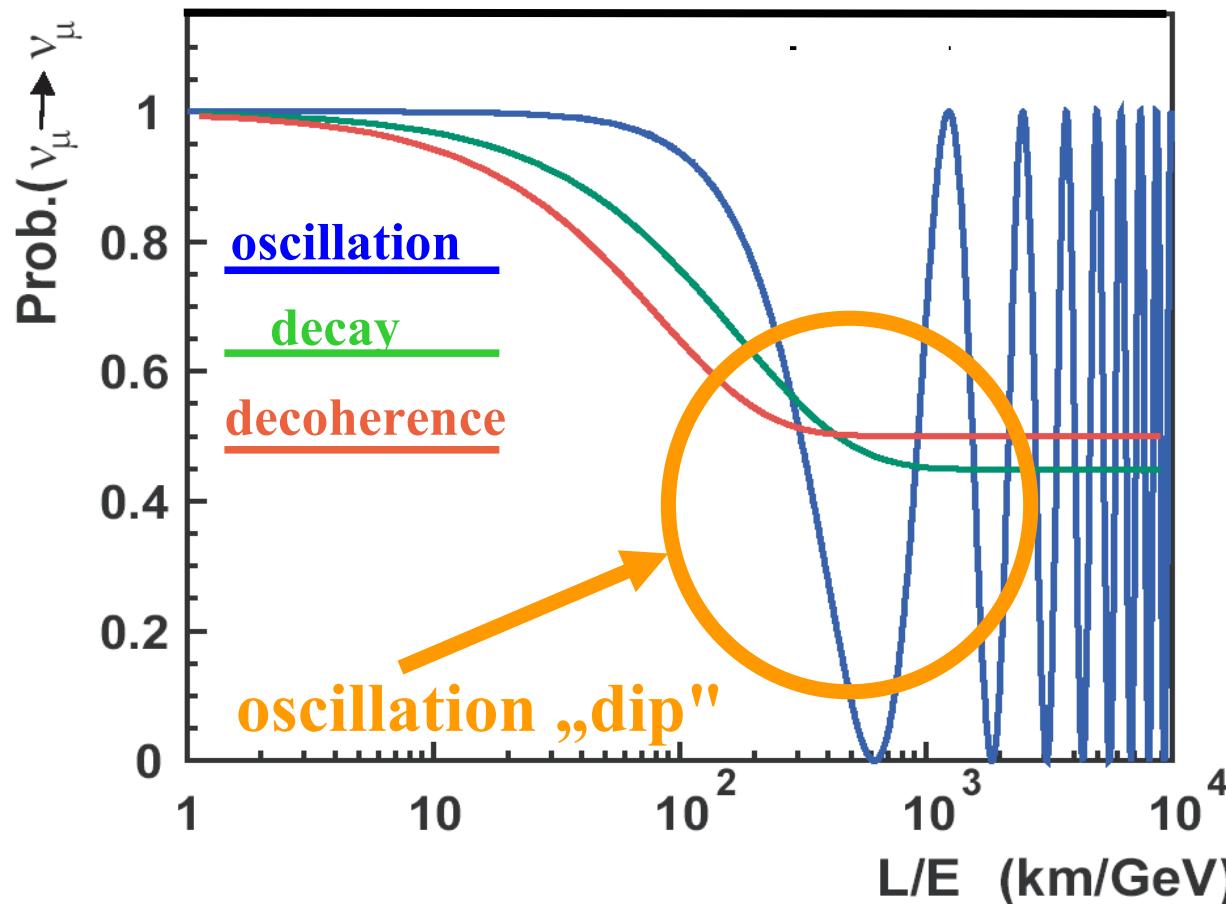
- 8σ signal for ν_μ disappearance
- NOT $\nu_\mu \rightarrow \nu_e$ (consistent with Chooz)
- $\Rightarrow \nu_\mu \rightarrow \nu_\tau$ (some τ' s seen)
- NOT $\nu_\mu \rightarrow \nu_s$ from NC/CC comparison
- L/E confirmed by K2K ($\simeq 2\sigma$)
- sensitivity for L/E-dependence of oscillations

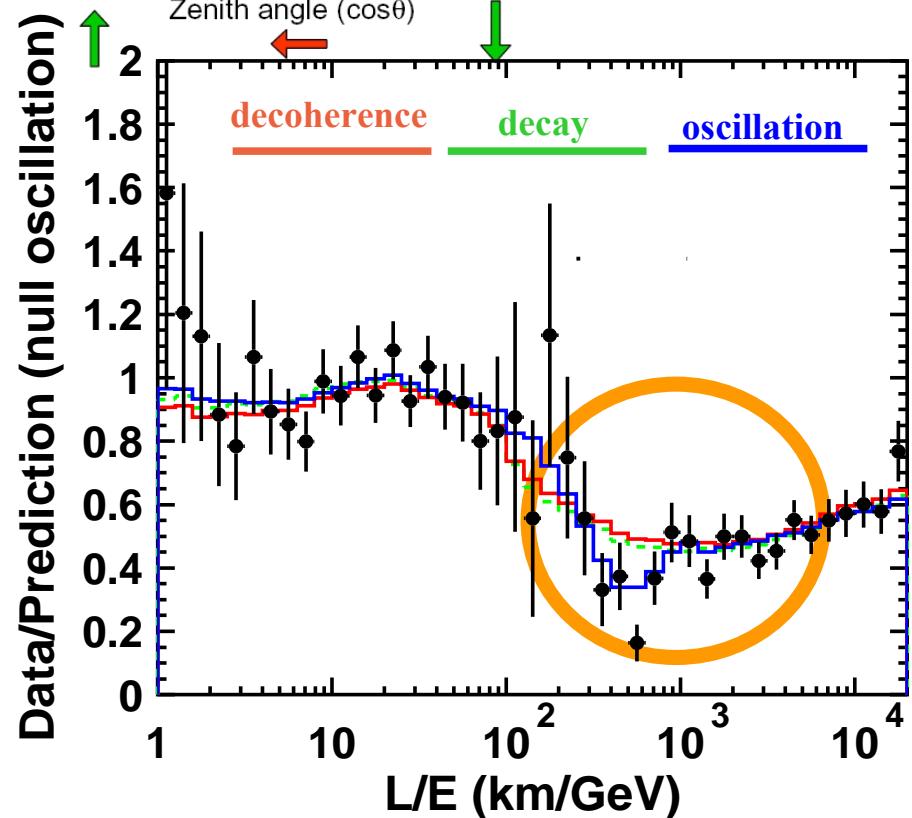
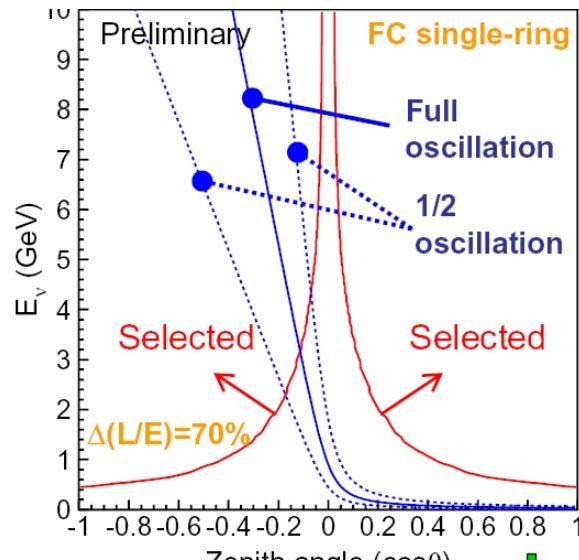
L/E Dependence: Atmospheric Oscillations

Neutrino oscillation : $P_{\mu\mu} = 1 - \sin^2 2\theta \sin^2(1.27 \frac{\Delta m^2 L}{E})$

Neutrino decay : $P_{\mu\mu} = (\cos^2 \theta + \sin^2 \theta \times \exp(-\frac{m}{2\tau} \frac{L}{E}))^2$

Neutrino decoherence : $P_{\mu\mu} = 1 - \frac{1}{2} \sin^2 2\theta \times (1 - \exp(-\gamma_0 \frac{L}{E}))$





Bad L/E resolution for:

- horizontal events ($dL/d\cos\theta$ is big!)
- events with small energy

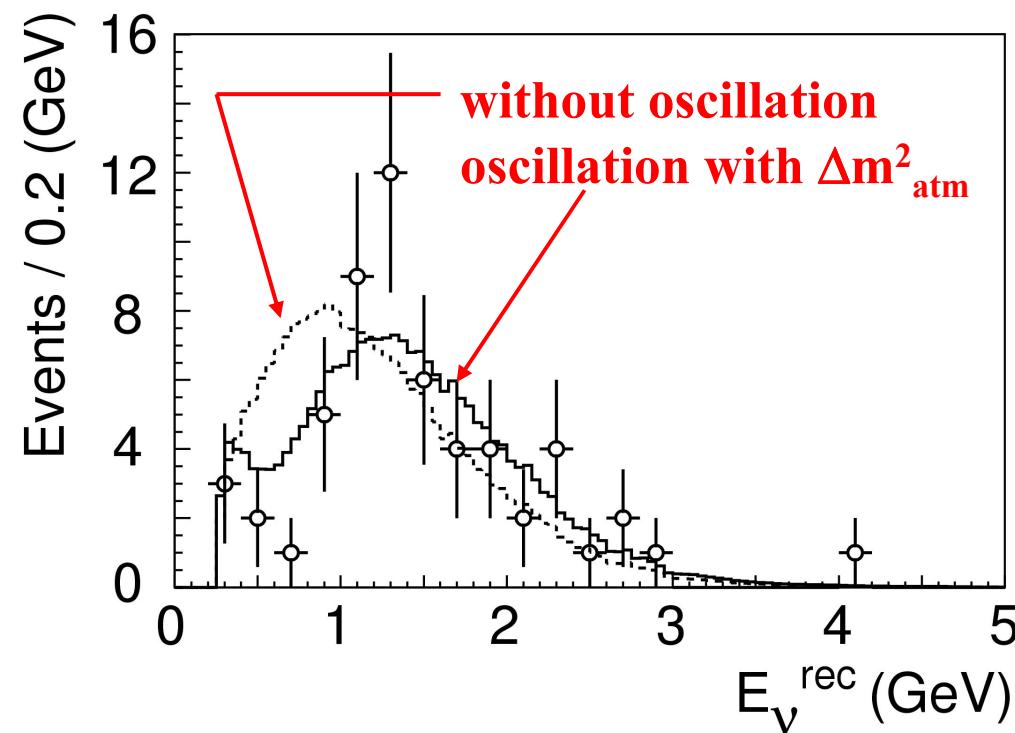
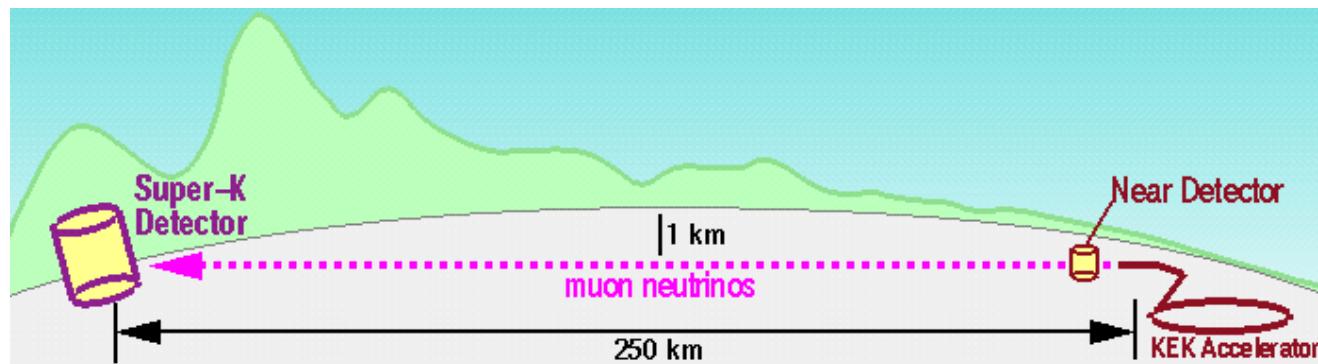
← cuts in the E-cosθ plane

SK I + II (preliminary)

Result:

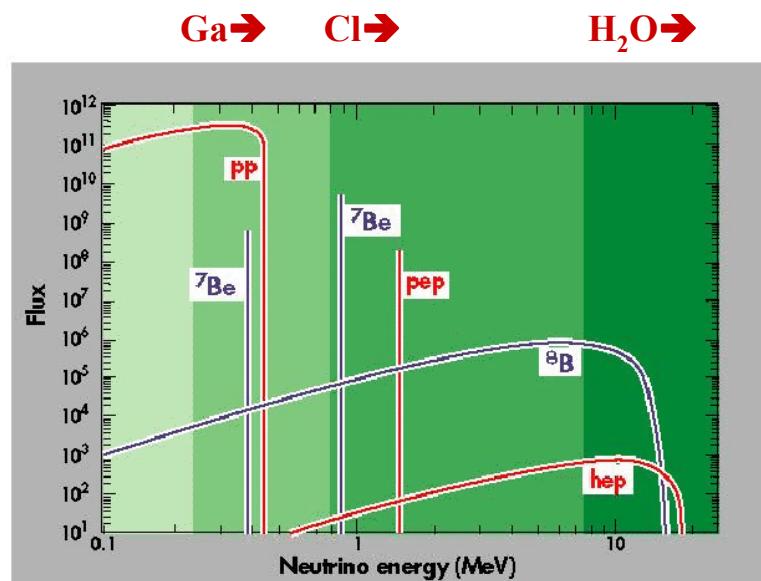
- 4.4σ for decay
- 4.8σ for de-coherence
- $\Delta m^2 = 2.4 \cdot 10^{-3} \text{ eV}^2$
- ←→ long baseline exp.

K2K confirms atmospheric Δm^2

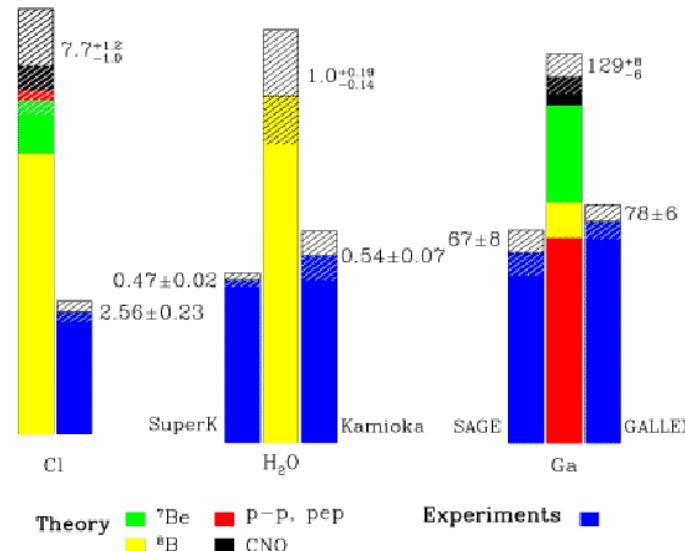


First Indications from Solar Neutrinos

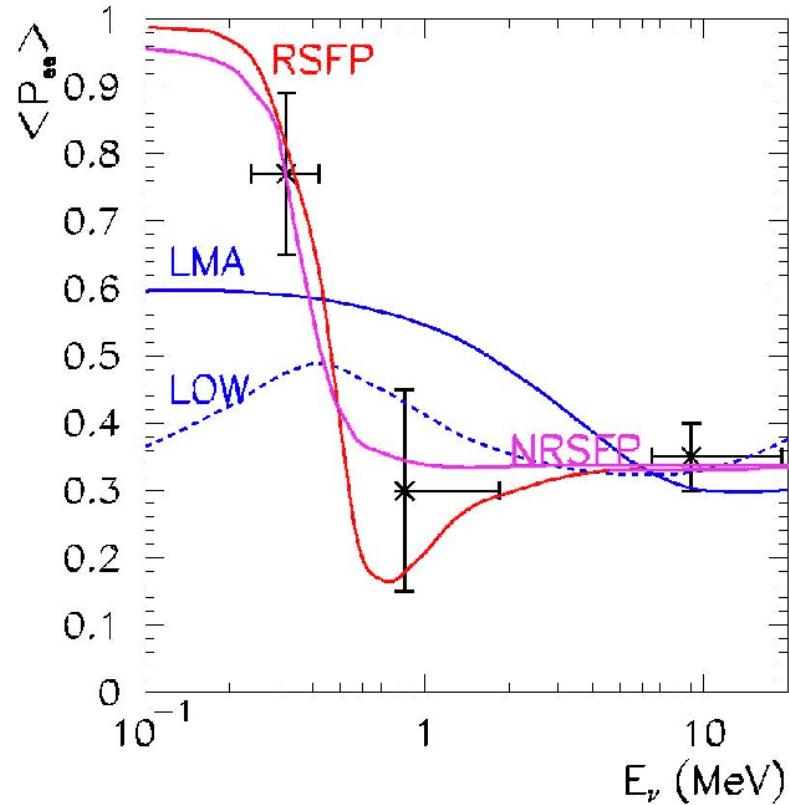
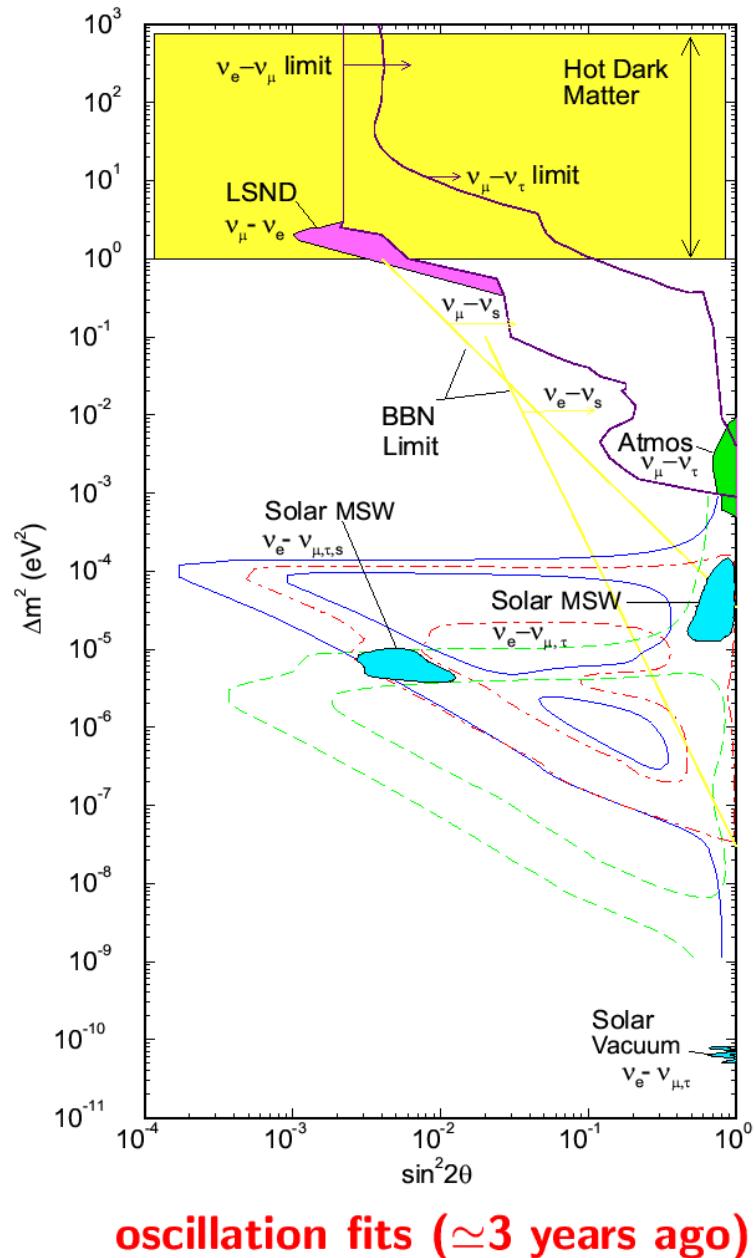
Before ~2000



Total Rates: Standard Model vs. Experiment
Bahcall-Pinsonneault 98

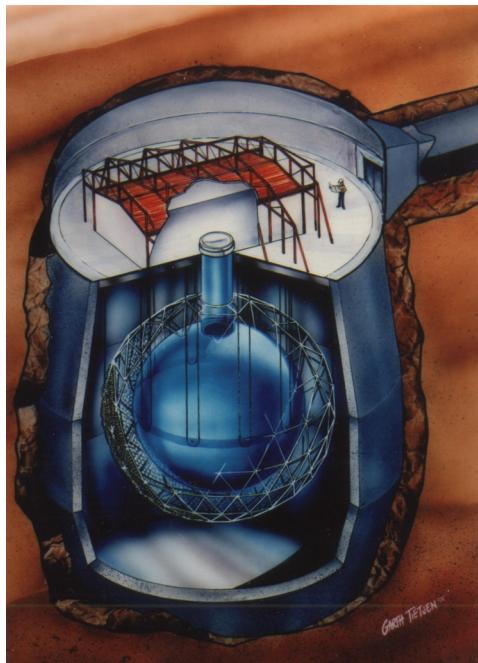


- Consistent deficit of total CC neutrino rates
- Combining *Cl/H₂O/Ga/D₂O* ⇒ some spectral information
- Depends on solar modelling, but should be robust
- Can be explained by oscillation ⇒ different regimes ⇒



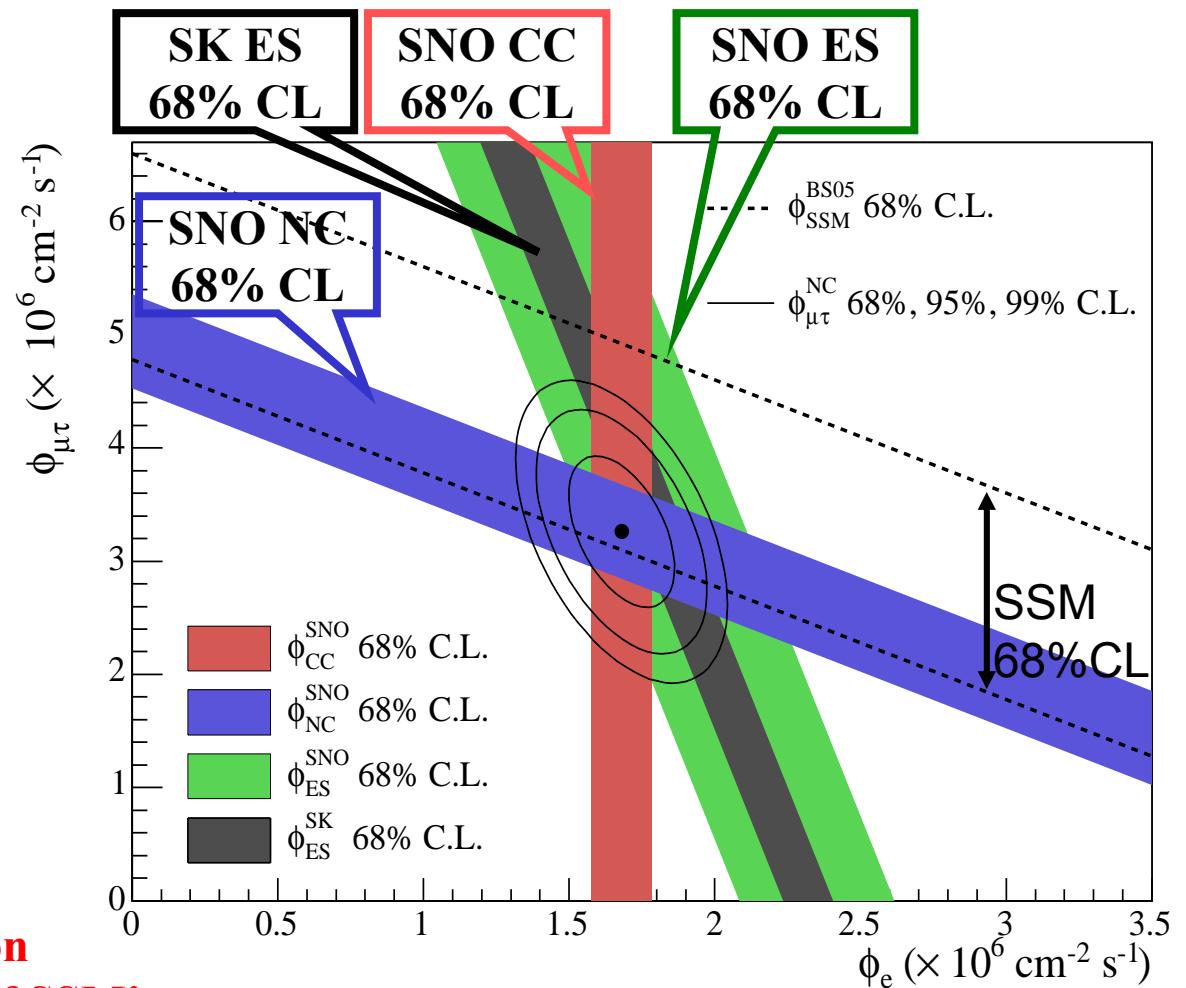
- **Which oscillation solution?**
 - SNO day/night \Rightarrow LMA best
- **RSFP worked even better!**
 - Requires large magnetic moments
 - ... and large magnetic fields

Solar ν 's: NC, CC, ES Rates from SNO



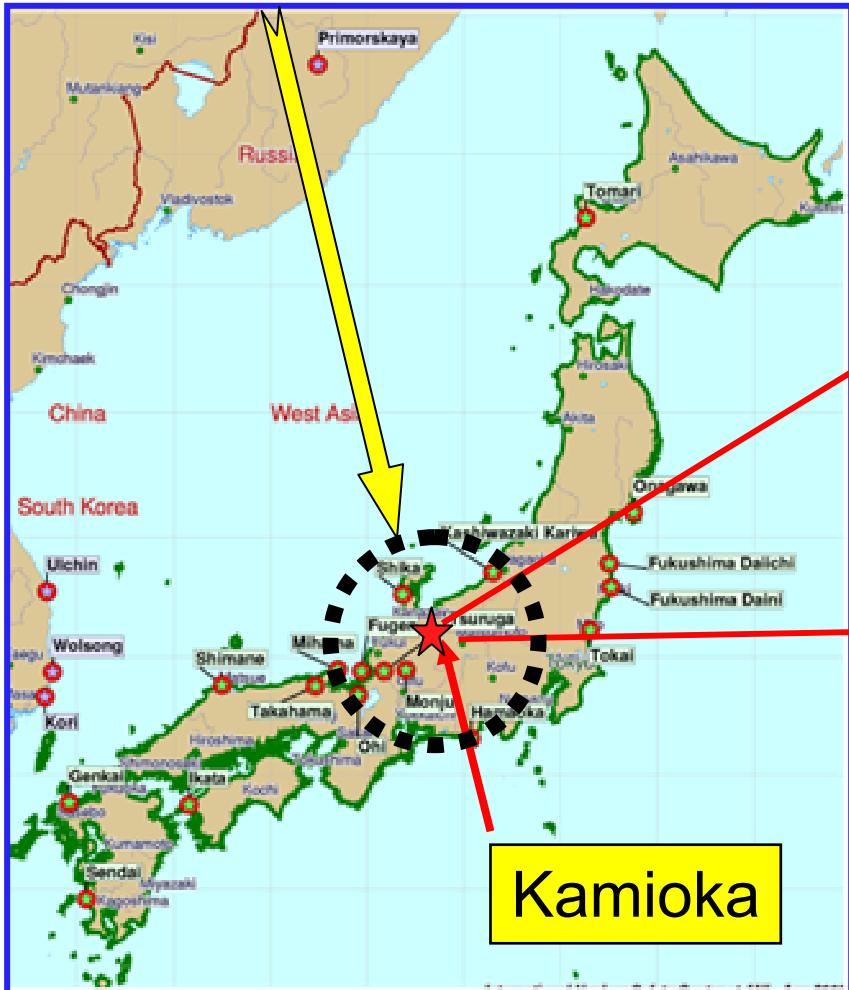
$$\begin{aligned}\Phi_{CC} &= \Phi_e \\ \Phi_{ES} &= \Phi_e + 1/6 \Phi_{\mu\tau} \\ \Phi_{NC} &= \Phi_e + \Phi_\mu + \Phi_\tau\end{aligned}$$

Clear proof of flavour conversion
Independent of SSM \leftrightarrow test of SSM!
No conversion and conversion to sterile ν 's excluded by many sigma
 \simeq no L/E dependence

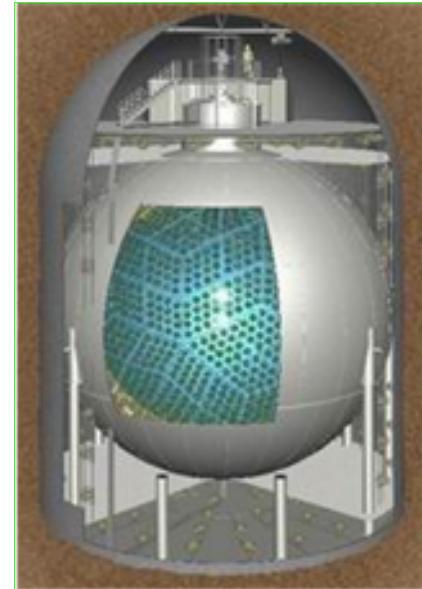


KamLAND tests solar Oscillations

26 power reactors are located in narrow band at $L \sim 180\text{km}$ from KamLAND, producing 80GW_{th} , 7% of World reactor power !!



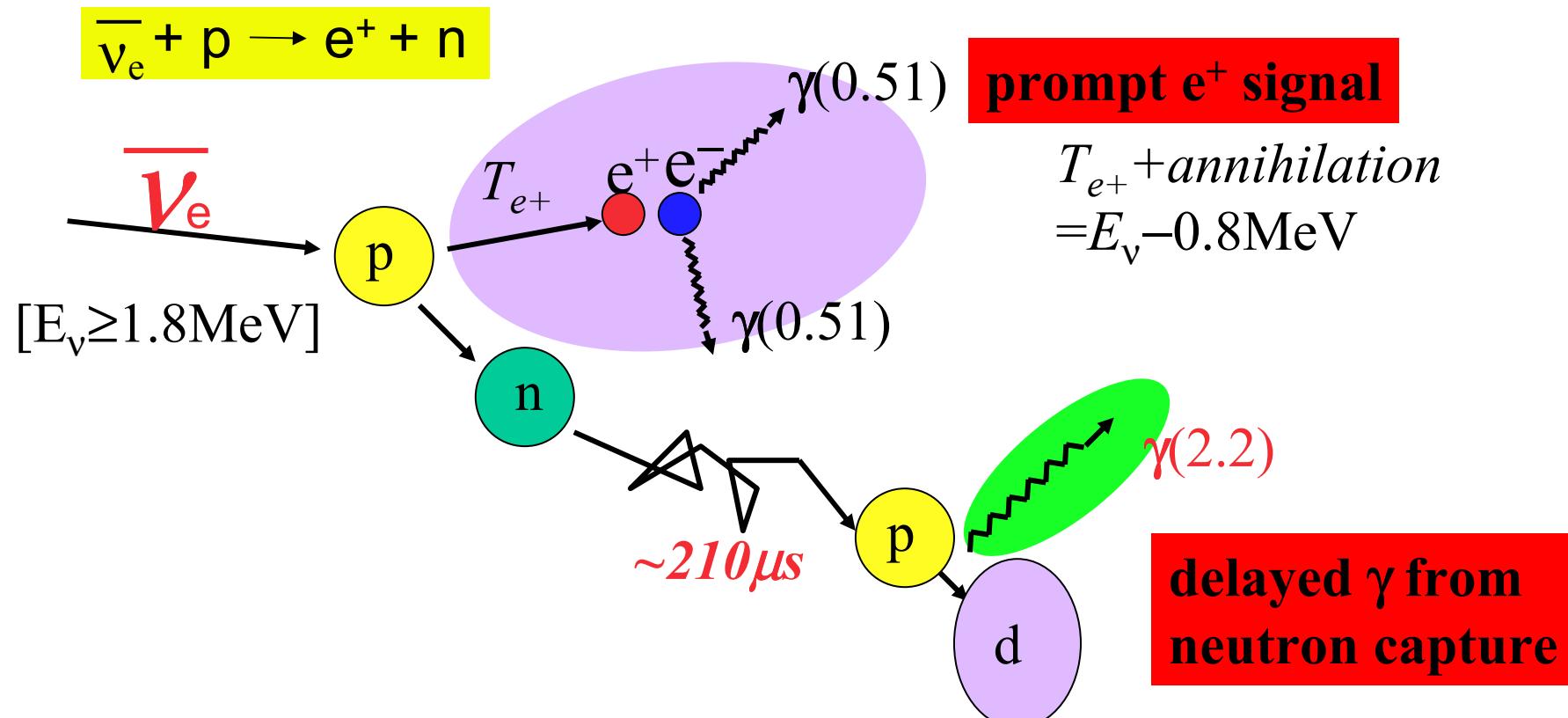
KamLAND



$\langle L \rangle \sim 180\text{km}$

Contribution from overseas
Korea 2.46%
Other countries 0.7%

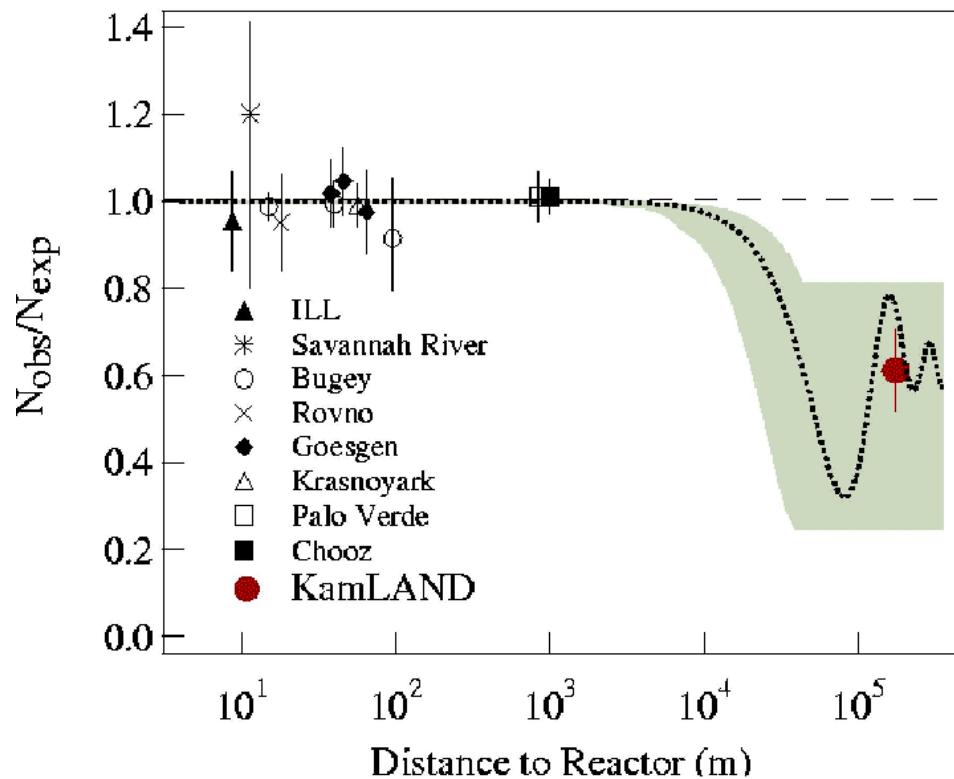
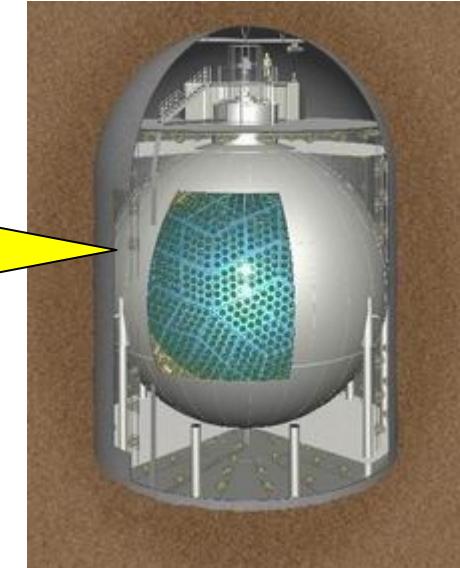
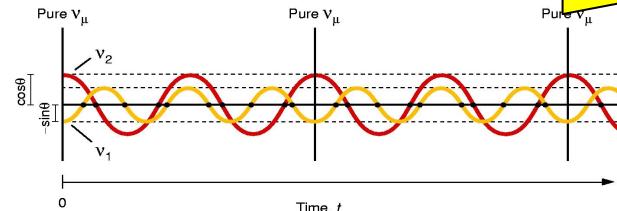
Neutrino Detection at KamLAND



position & time correlation
and delayed energy information
→ enormous background reduction!



flux $\sim 1/L^2$ + oscillations?



Compare to solar oscillations:

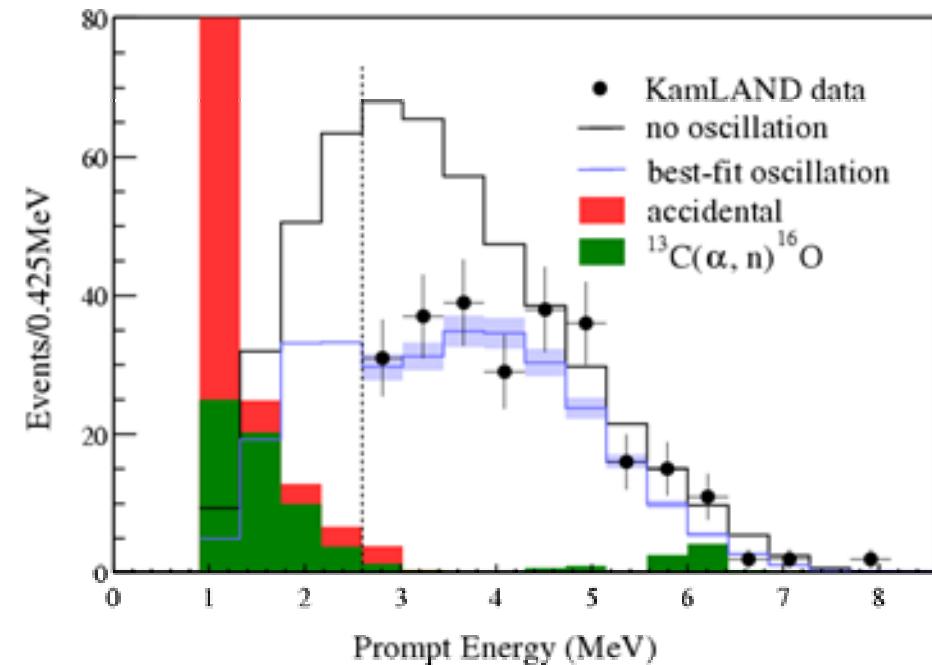
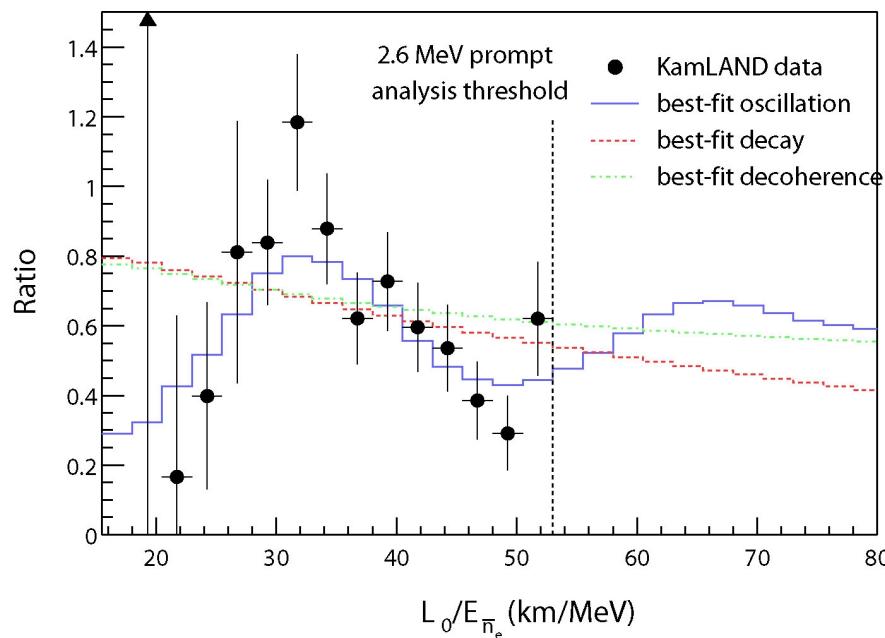
- rate consistent
- tests solar oscillations with reactor anti- ν 's

- excludes RSFP (no B fields)
- \simeq excludes CPT-violation \leftrightarrow LSND
- rate effect & some L/E sensitivity

KamLAND Results

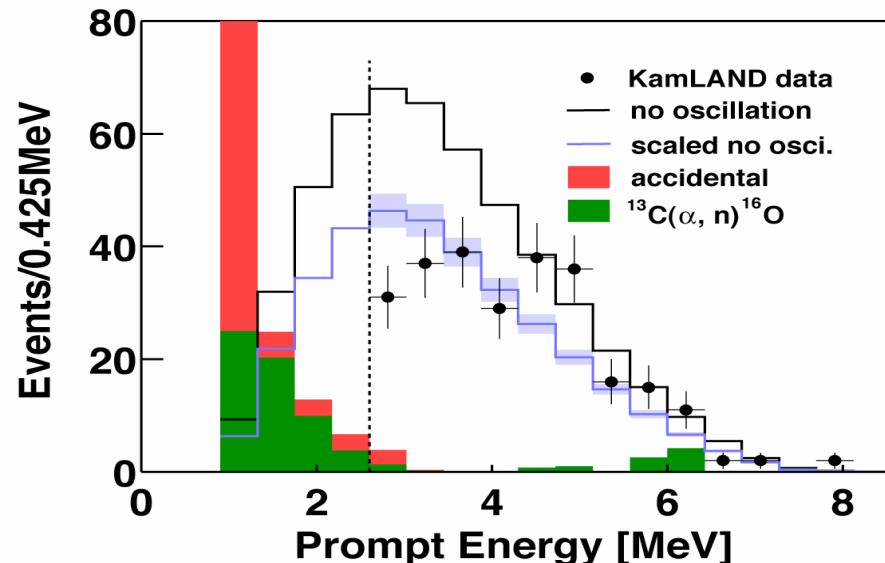
Best fit to latest results (2005):

$$\Delta m^2 = 7.9^{+0.6}_{-0.5} \times 10^{-5} \text{ eV}^2$$

$$\tan^2 \theta = 0.46$$


some L/E dependence

Testing Solar L/E with KamLAND

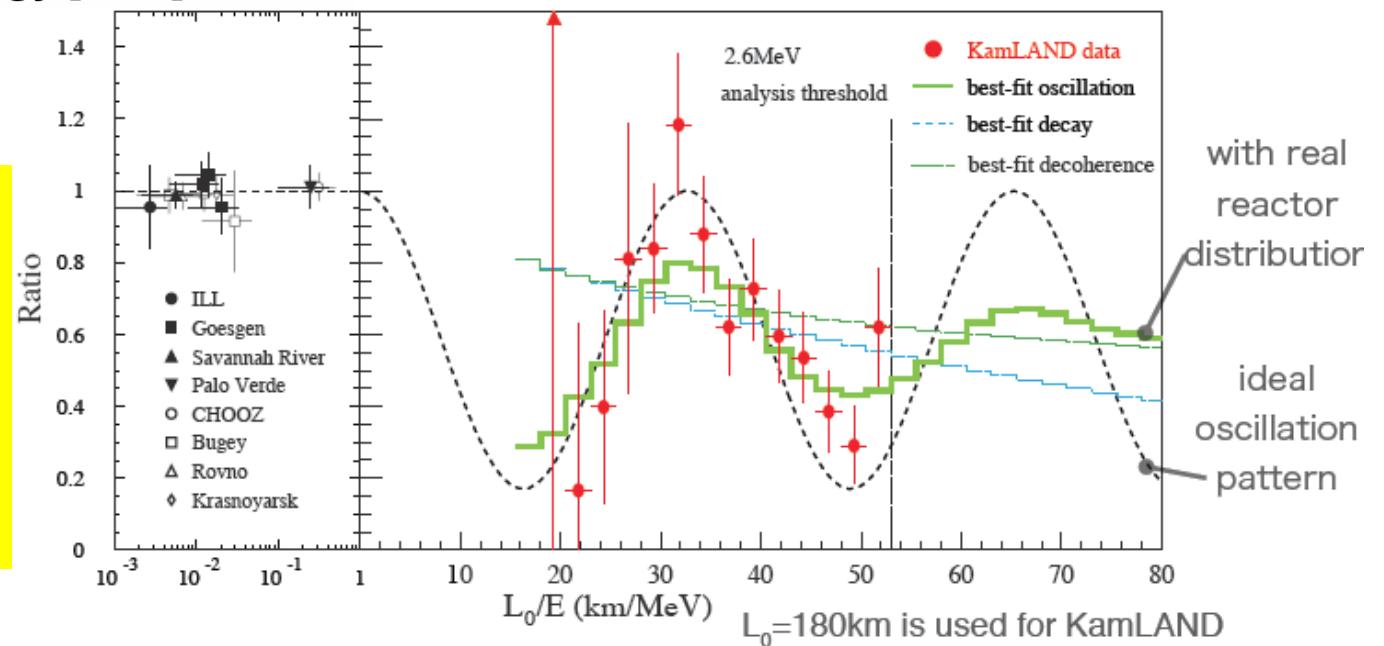


rate plus shape →
oscillations at 99.999995% CL

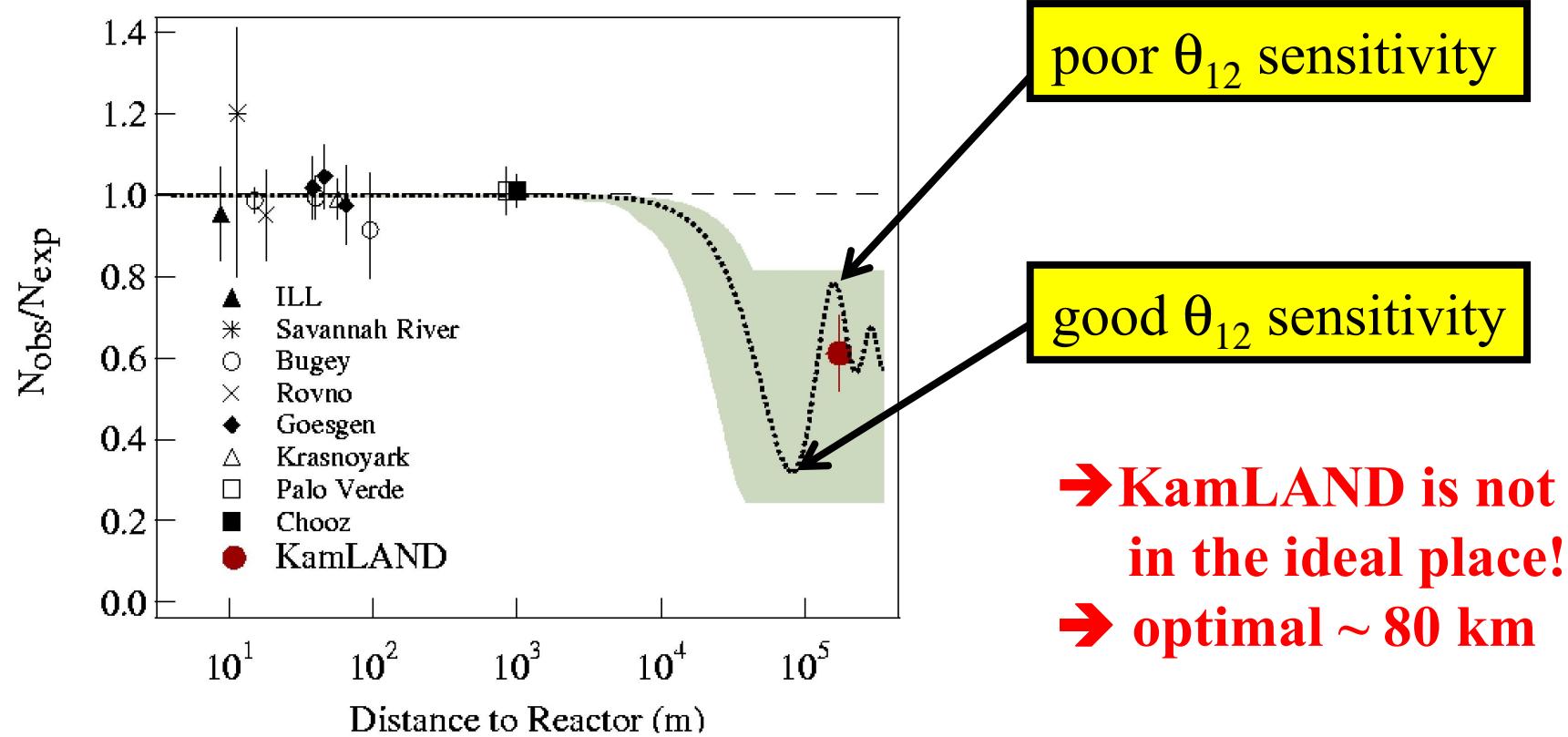
Best fit: $\Delta m^2 = 7.9^{+0.6}_{-0.5} \times 10^{-5} \text{ eV}^2$
 $\tan^2 \theta = 0.46$

improved tests of L/E:

- Super Kamiokande
- KamLAND
- MINOS
- ...

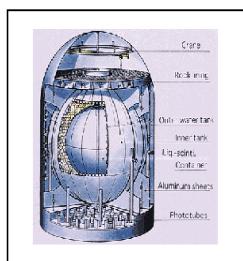


KamLAND Implications



- New reactor experiments at ideal location(s)
- better θ_{12} sensitivity
 - better Δm^2 sensitivity

Neutrino Oscillations Status Summary



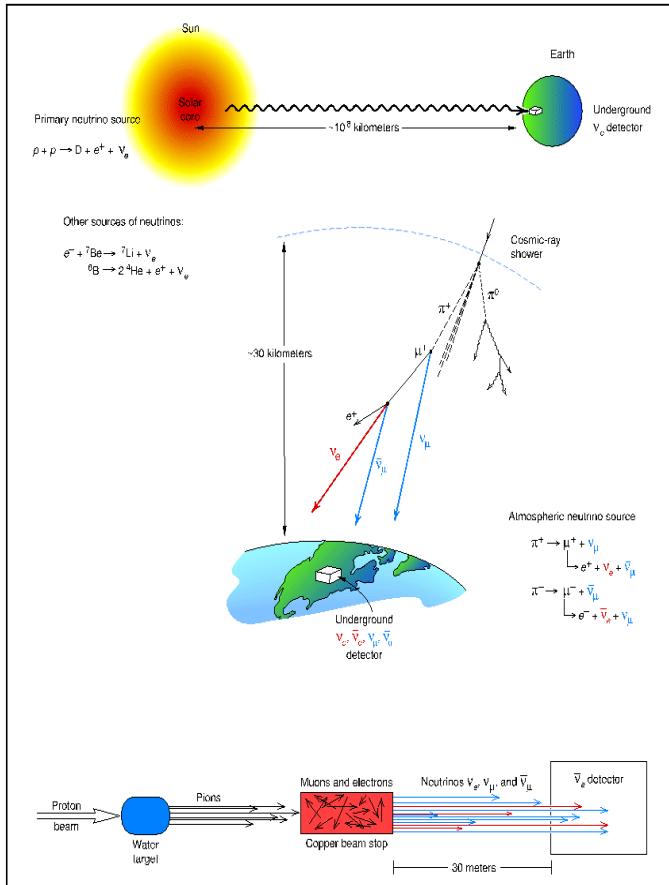
KamLAND
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$\sin^2 \theta_{13} \leq 0.041 @ 3\sigma$

Maltoni, Schwetz, Tortola, Valle

